

that the systems would differ considerably. Lynds, *et al.*,¹¹ presented two equations to describe the decrease in the lattice parameter as the oxide content of the lattice increases: one for the composition range UO_2 – $\text{UO}_{2.125}$ and another for the range $\text{UO}_{2.17}$ – $\text{UO}_{2.25}$. The first equation characterizes the behavior of a nonstoichiometric UO_2 phase while the second characterizes that of a U_4O_9 phase ($\text{UO}_{2.25}$). The first of these has a clear counterpart in the samarium fluoride system but the second does not, probably because a clearly defined $\text{SmF}_{2.25}$ analog could not be found. However, a single-crystal study by Belbeoch, *et al.*,¹² has demonstrated that the unit cell of $\text{UO}_{2.25}$ contains 64 nominally fluorite unit cells arranged in a $4 \times 4 \times 4$ array. The superstructure observed in the latter half of the cubic–tetragonal two-phase region and assigned to the tetragonal unit cell may, in fact, correspond to a similar type of cell and an $\text{SmF}_{2.25}$ phase may exist. The concurrence of a phase boundary at $\text{SmF}_{2.16}$ and $\text{UO}_{2.17}$ provides some support for the existence of a phase at $\text{SmF}_{2.25}$.

Another phase in the uranium–oxygen system, $\text{UO}_{2.33}$ (U_6O_{14}), a member of a homologous series ($\text{U}_n\text{O}_{2n+2}$) reported for this system by Makarov,¹⁰ has a similar analog in the samarium system, $\text{SmF}_{2.35}$. Hoekstra, *et al.*,¹³ have reported that this uranium phase possesses body-centered tetragonal symmetry with $a = 3.78$ and $c = 5.55$ Å. As in the samarium phase, the a parameter is less than half the face diagonal of a cube

(12) B. Belbeoch, C. Piekarski, and P. Perio, *Acta Cryst.*, **14**, 837 (1961).

(13) H. R. Hoekstra, A. Santoro, and S. Seigel, *J. Inorg. Nucl. Chem.*, **18**, 166 (1961).

with edge c . A modification of this tetragonal cell in which $a = 3.84$ and $c = 5.40$ Å has been observed for only the uranium–oxygen system.

The next phase in the homologous series, $\text{UO}_{2.40}$ (U_5O_{12}), has been reported to be tetragonal with $a = 5.364$ and $c = 5.531$ Å. The phase $\text{SmF}_{2.41}$, which has been indexed on a rhombohedral lattice, seems to be related to this uranium analog.

The samarium analog of $\text{UO}_{2.5}$ (U_4O_{10}) may correspond to $\text{SmF}_{2.46}$, but the relationship is not so well defined. The uranium compound has been examined by single-crystal techniques and has been found to be orthorhombic.¹⁰ It seems apparent that while many of the same stoichiometries are involved, the lattice modifications are different for the two systems. The uranium system displays at least one more phase, with several crystal modifications, before the lattice collapses to that of UO_3 , and thus appears to be more complicated at the anion-rich end of the nonstoichiometric region. This complexity probably results from the different charges and sizes of the ions.

It is apparent from the above discussion that the nonstoichiometry displayed by these two systems is the result of similar modifications of the fluorite lattice. The present investigation has been unable to establish the nature of these changes and has indicated that single-crystal studies must be undertaken to clarify some of the observations.

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The Crystal and Molecular Structure of Dinitrosylbis(triphenylphosphine)iridium Perchlorate, $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$

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The structure of dinitrosylbis(triphenylphosphine)iridium perchlorate, $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$, has been determined from three-dimensional X-ray data collected by counter methods. The compound crystallizes in the space group $\text{C}_{2h}^2\text{-C2}/c$ of the monoclinic system with four molecules in a unit cell of dimensions $a = 17.090$ (16) Å, $b = 12.413$ (13) Å, $c = 17.165$ (17) Å, and $\beta = 108.02$ (11)°. The observed and calculated densities are 1.69 (± 0.02) and 1.68 g/cm³, respectively. Least-squares refinement of the structure has led to a final value of the conventional R factor (on F) of 0.043 for the 1597 independent reflections having $F^2 > 3\sigma(F^2)$. The crystal structure consists of well-separated, discrete, monomeric ions. The coordination geometry around the iridium atom is only approximately tetrahedral; the N–Ir–N angle, 154.2 (7)°, is much larger than expected. Possible reasons for this distortion are discussed. The iridium–nitrogen–oxygen bond angle of 164 (1)° deviates considerably from linearity. The perchlorate anion is disordered about two positions and has been refined as a rigid group, with normal tetrahedral angles of 109° 28' and chlorine–oxygen bond lengths of 1.43 Å.

Introduction

Crystal structure studies of $[\text{IrX}(\text{CO})(\text{NO})(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{BF}_4]$ ($\text{X} = \text{Cl}$ or I)^{1–3} showed that in each

(1) D. J. Hodgson, N. C. Payne, J. A. McGinness, R. G. Pearson, and J. A. Ibers, *J. Am. Chem. Soc.*, **90**, 4486 (1968).

(2) D. J. Hodgson and J. A. Ibers, *Inorg. Chem.*, **7**, 2345 (1968).

(3) D. J. Hodgson and J. A. Ibers, *ibid.*, **8**, 1282 (1969).

compound the metal–nitrosyl geometry was very different from that previously observed in reliable structural studies of transition metal–nitrosyl complexes. Whereas the usual transition metal–nitrosyl geometry involves a short metal–nitrogen bond and a linear or nearly linear metal–nitrogen–oxygen bond angle,

the structural studies of the iridium systems showed that the geometry involved a long metal-nitrogen bond and a metal-nitrogen-oxygen bond angle of approximately 120° . These studies led to the suggestion¹ that while in the usual transition metal-nitrosyl complexes the nitrosyl ligand is acting as a Lewis base (NO^+), in the iridium systems it is acting as a Lewis acid (NO^-). We were thus prompted to examine other nitrosyl complexes where the nitrosyl ligand might be acting as a Lewis acid.

Angoletta, Malatesta, and Caglio^{4,5} isolated $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$ by treating an ethanolic solution of $[\text{IrH}_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$ with nitric oxide. On the basis of its infrared spectrum, the complex was formulated as containing iridium(-I), with the nitrosyls coordinated as " NO^+ ." Such a d^{10} complex would be expected to be tetrahedral. However, the reactions of this complex are not completely consistent with this formulation. For example, the complex reacts with lithium halides to give the neutral five-coordinate compounds $\text{IrX}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2$ ($\text{X} = \text{Cl}, \text{Br}, \text{or I}$), which have two electrons greater than that predicted by the inert-gas rule. In such a complex the additional electrons probably reside on the amphoteric nitrosyl ligands rather than in antibonding metal-ligand σ orbitals. No similar reactions have been reported for the analogous complexes $\text{Pt}(\text{CO})_x(\text{P}(\text{C}_6\text{H}_5)_3)_{4-x}$ ($x = 0-2$).⁶ This unusual reaction may be rationalized in two ways. Either the starting complex is not iridium(-I) but iridium(I), or nucleophilic attack by the halide ions causes internal electron transfer to occur from one of the nitrosyl σ -bonding molecular orbitals to antibonding orbitals localized mainly on the nitrosyls thus leaving an empty σ -bonding orbital available for the lone pair of electrons of the halide ion. The iridium(I) ion may formally be written as $\text{Ir}(\text{NO}^+)(\text{NO}^-)(\text{P}(\text{C}_6\text{H}_5)_3)_2^+$, although this is an oversimplification and the electrons may actually be delocalized over both nitrosyls. Such an iridium(I) ion would be expected to be square planar. The present structural study of $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$ was undertaken to gain an insight into the oxidation state of the iridium atom and perhaps to understand the unusual reactions of this complex.

Collection and Reduction of Intensity Data

Violet, well-formed crystals of $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$ were prepared by the method of Angoletta and Caglio⁴ and were characterized by their melting point, infrared spectrum, and elemental analysis. On the basis of Weissenberg and precession photography of the $h0l$, $h1l$, $h2l$, $0kl$, $1kl$, and $2kl$ nets using Mo $K\alpha$ and Cu $K\alpha$ radiation, it was established that the crystals belong to the monoclinic system. The observed extinctions $h + k = 2n + 1$ for hkl and $l = 2n + 1$ for $h0l$ suggested the space groups C_{2h}^6 - $\text{C}2/c$ or C_s^4 - Cc . On the basis of optical goniometry, the

crystal faces were identified as belonging to the forms $\{\bar{1}11\}$ and $\{100\}$. The lattice constants at 23° , which were determined from a least-squares refinement of the setting angles of 21 strong X-ray reflections which had been centered on a Picker four-circle automatic diffractometer⁷ using Mo $K\alpha_1$ radiation (λ 0.70930 Å), are $a = 17.090$ (16) Å, $b = 12.413$ (13) Å, $c = 17.165$ (17) Å, and $\beta = 108.02$ (11) $^\circ$. The density calculated for four formula weights per unit cell is 1.68 g/cm³, which agrees well with that of 1.69 (± 0.02) g/cm³ measured by suspending crystals in a mixture of diiodomethane and carbon tetrachloride. With four formula weights per unit cell either a two-fold axis or a center of symmetry is imposed on the ions in the space group $\text{C}2/c$ but no symmetry is imposed in the noncentric space group Cc .

The mosaicity of the crystal was examined by means of the narrow source, open counter, ω -scan technique. The width at half-height for a typical strong reflection was found to be approximately 0.05° , which is acceptably low.⁸ The intensity data were collected as previously described⁷ using a crystal with approximate dimensions $0.53 \times 0.063 \times 0.063$ mm. The crystal was mounted on the diffractometer with the needle direction ($[101]$ zone axis) almost coincident with the diffractometer spindle axis.

For data collection Mo $K\alpha$ radiation was used and the diffracted beams were filtered through 3.0 mils of zirconium foil. The intensities were measured by the θ - 2θ technique at a takeoff angle of 1.3° . At this angle the intensity of a reflection was about 80% of the maximum value as a function of takeoff angle. A receiving counter aperture 4×4 mm was used and was positioned 35 cm from the crystal. A nonsymmetric scan range in 2θ was used, from -0.5 to $+1.0$ from the $K\alpha_1$ peak. Stationary counter-stationary crystal background counts of 10 sec were measured at each end of the scan. Attenuators were inserted automatically if the intensity of the diffracted beam exceeded about 7000 counts/sec during the scan; the attenuators used were Cu foils, whose thickness had been chosen to give attenuator factors of approximately 2.3. The pulse height analyzer was set for approximately a 90% window, centered on the Mo $K\alpha$ peak.

The unique data set having $2\theta < 40^\circ$ was gathered; the intensities of a total of 1882 independent reflections were recorded. The intensities of five standard reflections, measured after every 200 reflections, remained essentially constant throughout the run, showing only the deviations of less than $\pm 1\%$. There were very few reflections above background at values of $2\theta > 40^\circ$.

All data processing was carried out as previously described.⁷ The value of p was selected as 0.04. The values of I and $\sigma(I)$ were corrected for Lorentz and polarization effects. Of the 1882 reflections, 144 had

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(5) M. Angoletta, *Gazz. Chim. Ital.*, **93**, 1591 (1963).

(6) R. Ugo, *Coord. Chem. Rev.*, **3**, 319 (1968).

(7) P. W. R. Corfield, R. J. Doedens, and J. A. Ibers, *Inorg. Chem.*, **6**, 197 (1967).

(8) T. C. Furnas, "Single Crystal Orienter Instruction Manual," General Electric Co., Milwaukee, Wis., 1957.

intensities which were less than their standard deviations. The linear absorption coefficient, μ , for this compound using Mo K α radiation is 43.30 cm⁻¹. For this crystal trial calculations showed that the transmission coefficients ranged only from 0.78 to 0.75 and therefore no absorption correction was made.

Solution and Refinement

The space group was assumed to be C2/c and the iridium and phosphorus atoms were located from a three-dimensional, origin-removed, sharpened Patterson function.⁹ The position of the iridium at 0, y , $1/4$ requires that twofold symmetry be imposed on the cation. A least-squares refinement on F was computed and the function $\sum w(|F_o| - |F_c|)^2$ was minimized, where $w = 4F_o^2/\sigma^2(F_o^2)$, and $|F_o|$ and $|F_c|$ are the observed and calculated structure amplitudes. In all calculations of F_c , the atomic scattering factors of iridium were taken from Cromer and Waber,¹⁰ and those for hydrogen, from Stewart, Davidson, and Simpson.¹¹ Scattering factors for all other atoms were taken from the usual tabulation.¹² The effects of anomalous dispersion were included in the calculations of F_c .¹³ The values of $\Delta f'$ and $\Delta f''$ were taken from Cromer's calculations.¹⁴ Only the 1597 reflections which had intensities greater than three times their standard deviations were included in the structural refinement. The iridium and phosphorus atoms were assigned isotropic thermal parameters and one cycle of least-squares refinement gave the agreement factors $R_1 = 0.229$ and $R_2 = 0.302$, where $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$ and R_2 (weighted R factor) = $(\sum w(|F_o| - |F_c|)^2 / \sum w F_o^2)^{1/2}$.

A subsequent difference Fourier synthesis led to positions of all the nonhydrogen atoms of the cation and to that of the perchlorate chlorine atom. The parameters obtained from this map were refined by least-squares methods. The nonphenyl atoms were refined with isotropic thermal parameters; the phenyl rings were treated as rigid groups and restricted to their known geometry (D_{6h} symmetry, $d(C-C) = 1.390$ Å). The variable parameters for the rings included an overall isotropic thermal parameter, the coordinates of the ring center, and three orientation angles, which have been defined previously.¹⁵⁻¹⁷ Two cycles of refinement gave the agreement factors $R_1 = 0.083$ and $R_2 = 0.109$. Examination of another difference Fourier map revealed 14 of the 15 phenyl hydrogen atoms. This in part justifies our original

assumption that the compound belongs to the space group C2/c rather than Cc.

Although the perchlorate group could possess a crystallographic twofold axis and be ordered, examination of the difference Fourier map revealed that this was not so but rather that the perchlorate group was disordered over two sites related by the twofold axis. Three of the four independent half-oxygen atoms displayed peak heights of 2.2–2.8 e⁻/Å³ while the fourth half-oxygen atom had a peak height of 1.3 e⁻/Å³. The geometry, derived from the interpolated positions of the peak centers, was a reasonable one for a perchlorate group.

At this point the phenyl proton contributions to the structure factors were included as a fixed contribution ($d(C-H) = 1.08$ Å and the isotropic group thermal parameters were made equal to the corresponding carbon group thermal parameters). In a further refinement cycle all nongroup atoms, except for the perchlorate oxygens, were allowed to vibrate anisotropically and individual, variable isotropic thermal parameters were assigned to each of the ring carbon atoms. Two cycles of refinement yielded agreement factors of $R_1 = 0.040$ and $R_2 = 0.049$ for the 90 variables and 1597 observations. However, this treatment yielded an unsatisfactory geometry for the perchlorate group. Although the O–Cl–O bond angles were within the tolerable range of 99–121°, three of the Cl–O bond lengths were too long (1.50–1.56 Å) and the fourth was too short (1.20 Å). Such a disparity in bond lengths could be the result of disorder of the chlorine atom about the twofold axis. In the final round of calculations the perchlorate oxygen atoms were treated as a group with Cl–O distances of 1.43 Å and O–Cl–O angles of 109° 28' and were assigned individual isotropic temperature factors. The origin of the O₄ group was restricted to the twofold axis. After three cycles of least-squares refinement, the agreement factors were $R_1 = 0.043$ and $R_2 = 0.052$ (with 82 variables). The final parameters for the cation are essentially the same as those calculated for the first model. A final difference Fourier map revealed two peaks of heights 1.2 and 1.0 e⁻/Å³ at 1.5 and 1.6 Å from the chlorine atom. Presumably our model for the disorder is not completely satisfactory.

Two other models for the disorder were examined, but neither was as satisfactory as the one above. One model treated the perchlorate oxygens as a group not restricted to the twofold axis and the chlorine atom was treated as an individual atom with anisotropic temperature factors. This calculation led to agreement factors similar to those above but gave unreasonable Cl–O bond lengths as the centroid of the group did not coincide with the position of the chlorine atom. The other model treated all five of the ClO₄ atoms as a group, and gave similar agreement factors. However a difference Fourier map revealed a considerable amount of electron density around the chlorine atom, probably because the chlorine atom was not refined anisotropically.

(9) In addition to various local programs, the programs used in the refinement of this structure were local modifications of Hamilton's GONO9 absorption correction program, Busing and Levy's ORFFE function and error program, Johnson's ORTEP thermal ellipsoid plotting program, and Zalkin's FORDAP Fourier program. Our local least-squares program NUCLS in its nongroup form closely resembles the Busing–Levy ORFLS program.

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(12) J. A. Ibers in "International Tables of X-Ray Crystallography," Vol. 3, The Kynoch Press, Birmingham, England, Table 3.3.1.A.

(13) J. A. Ibers and W. C. Hamilton, *Acta Cryst.*, **17**, 781 (1964).

(14) D. T. Cromer, *ibid.*, **18**, 17 (1965).

(15) S. J. La Placa and J. A. Ibers, *J. Am. Chem. Soc.*, **87**, 2581 (1965).

(16) S. J. La Placa and J. A. Ibers, *Inorg. Chem.*, **5**, 405 (1966).

(17) R. Eisenberg and J. A. Ibers, *ibid.*, **4**, 773 (1965).

TABLE I
 POSITIONAL, THERMAL, AND GROUP PARAMETERS FOR $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$

Atom	<i>x</i>	<i>y</i>	<i>z</i>	β_{11}^a	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Ir	0	0.03051 (5) ^b	1/4	0.00300 (3)	0.00495 (5)	0.00320 (3)	0	0.00335 (2)	0
Cl	1/2	0.1154 (5)	1/4	0.0061 (3)	0.0158 (6)	0.0036 (2)	0	0.0016 (2)	0
P	0.0994 (2)	0.1299 (2)	0.3480 (2)	0.0028 (1)	0.0051 (2)	0.0028 (1)	0.0000 (1)	0.0010 (1)	-0.0002 (1)
N	0.0552 (6)	-0.0014 (8)	0.1811 (7)	0.0041 (5)	0.0097 (10)	0.0045 (5)	0.0011 (5)	0.0000 (4)	-0.0019 (5)
O	0.0923 (7)	0.0046 (10)	0.1314 (7)	0.0064 (5)	0.0231 (15)	0.0059 (5)	0.0022 (7)	0.0024 (5)	-0.0038 (7)
Group	<i>x</i> _c ^c	<i>y</i> _c	<i>z</i> _c	δ	ϵ	η			
R1	0.2065 (3)	0.2827 (4)	0.2700 (3)	0.407 (5)	-2.660 (5)	-0.765 (5)			
R2	0.2279 (3)	-0.0243 (4)	0.4746 (3)	2.294 (4)	-2.958 (5)	2.597 (5)			
R3	0.0266 (3)	0.2723 (4)	0.4648 (3)	0.942 (6)	2.434 (4)	1.333 (6)			
O4 ^d	1/2	0.1157 (13)	1/4	-2.06 (4)	-1.86 (1)	-0.70 (4)			

^a The form of the anisotropic ellipsoid is $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. ^b Standard deviations of the least significant figures are given in parentheses here and in subsequent tables. ^c *x*_c, *y*_c, and *z*_c are the fractional coordinates of the rigid group centers. The angles δ , ϵ , and η (in radians) have been defined previously.¹⁵⁻¹⁷ ^d O4 represents the tetrahedral group of the perchlorate oxygen atoms.

TABLE II

Atom *x* *y* *z* *B*, Å²a. Derived Parameters for Ring Carbon Atoms^a

R1C1	0.1607 (5)	0.2177 (6)	0.3056 (4)	2.8 (2)
R1C2	0.2096 (5)	0.1712 (4)	0.2633 (5)	4.4 (2)
R1C3	0.2554 (4)	0.2362 (7)	0.2277 (4)	5.3 (3)
R1C4	0.2524 (6)	0.3476 (6)	0.2344 (5)	6.0 (3)
R1C5	0.2035 (6)	0.3941 (4)	0.2767 (5)	6.3 (3)
R1C6	0.1577 (4)	0.3291 (6)	0.3123 (4)	4.4 (2)
R2C1	0.1734 (4)	0.0404 (5)	0.4165 (4)	3.3 (2)
R2C2	0.2553 (4)	0.0712 (4)	0.4496 (5)	4.0 (2)
R2C3	0.3098 (3)	0.0064 (6)	0.5076 (5)	4.6 (3)
R2C4	0.2824 (4)	-0.0891 (6)	0.5326 (4)	5.5 (3)
R2C5	0.2006 (5)	-0.1198 (5)	0.4996 (5)	5.5 (3)
R2C6	0.1461 (3)	-0.0551 (6)	0.4416 (5)	4.1 (2)
R3C1	0.0585 (6)	0.2094 (7)	0.4148 (4)	3.1 (2)
R3C2	0.0881 (5)	0.1976 (7)	0.4995 (4)	4.3 (2)
R3C3	0.0562 (5)	0.2605 (6)	0.5494 (3)	5.3 (3)
R3C4	-0.0053 (6)	0.3352 (7)	0.5148 (4)	5.3 (3)
R3C5	-0.0349 (5)	0.3471 (7)	0.4302 (5)	5.4 (3)
R3C6	-0.0030 (4)	0.2842 (5)	0.3802 (3)	4.2 (2)

b. Derived Parameters for Perchlorate Group Oxygen Atoms

O2	0.501 (3)	0.034 (1)	0.192 (1)	11.3 (7)
O3	0.473 (4)	0.215 (2)	0.208 (1)	11.1 (7)
O4	0.581 (1)	0.129 (5)	0.306 (1)	18.5 (13)
O5	0.445 (1)	0.085 (4)	0.294 (1)	9.9 (6)

^a Cl is attached to P; other C atoms are numbered in succession so that C4 is *para* to Cl.

All these calculations had only marginal effects on the positional parameters of the cation. The nitrogen and oxygen parameters were most affected. We feel that, in view of the disorder of the perchlorate group, the standard deviations of the bond lengths and angles associated with the nitrogen and oxygen atoms may be optimistic, but not by more than a factor of 2. The thermal parameters of the perchlorate oxygen atoms in the final model reflect the disparity in the peak heights found in the difference Fourier map. Why one oxygen atom should be vibrating more than the other three is not obvious from the mode of packing of the ions.

The value of *R*₂ obtained from the final cycle of least-squares refinement shows no dependence on $|F_o|$ or on $\lambda^{-1} \sin \theta$; thus the choice of *p* of 0.04 is essentially correct. Comparison of the final values of $|F_o|$ and $|F_c|$ suggests that no correction for secondary extinction

is necessary. The error of an observation of unit weight is 1.59 electrons.

The positional, thermal, and group parameters derived from the last cycle of least-squares refinement are given in Table I, along with the associated standard deviations as estimated from the inverse matrix. The positional parameters of the perchlorate group oxygen atoms and the ring carbon atoms which may be derived from the data in Table I are presented in Table IIa and b together with their thermal parameters as obtained from the last cycle of least-squares refinement. The final values of $5|F_o|$ and $5|F_c|$ (in electrons) are given in Table III; only the 1597 reflections which were used in the refinement are listed in this table. Of the 285 reflections omitted from the refinement for which $F_o^2 < 3\sigma(F_o^2)$ 284 had $|F_o^2 - F_c^2| < 3\sigma(F_o^2)$. The reflection 919 had $|F_o^2 - F_c^2| = 5.7\sigma(F_o^2)$.

Description of Structure

The crystal structure consists of discrete monomeric ions, with the cation on and the anion disordered about crystallographically imposed twofold axes of symmetry. The nearest contact between the iridium atom and the perchlorate oxygen fragments is 3.98 Å.

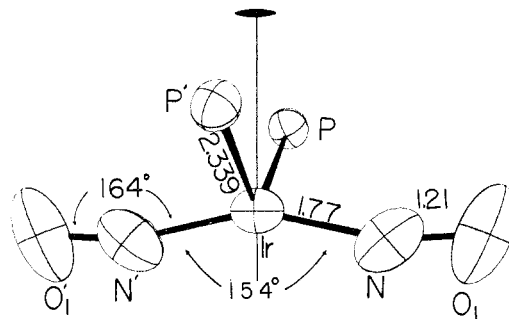


Figure 1.—The inner coordination sphere of the cation. Atoms related by the twofold axis are represented by X and X'.

The inner coordination geometry of the cation is shown in Figure 1 and a stereoscopic view of the cation is shown in Figure 2.

A selection of distances and angles is given in Table

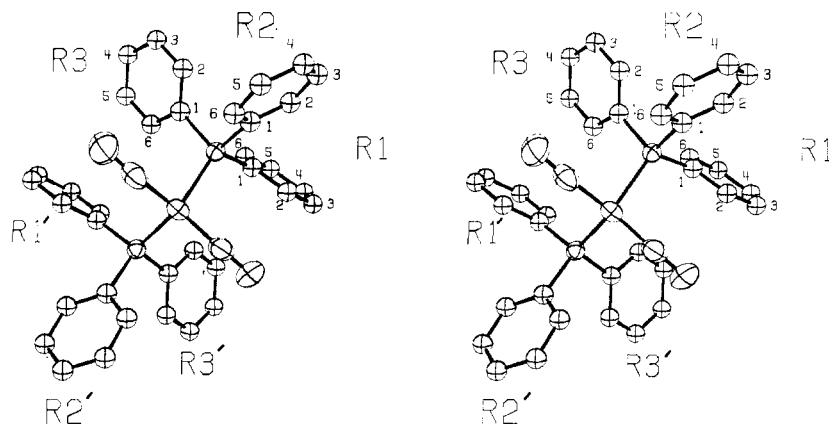


Figure 2.—A stereoscopic view of the cation $\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2^+$. The isotropic thermal parameters of the carbon atoms have all been set to 3.0 \AA^2 . For other atoms 50% probability ellipsoids are shown.

TABLE IV
SELECTED DISTANCES AND ANGLES^a

—Cation intramolecular dist, Å—		—Angle, deg—	
Ir—P	2.339 (3)	P—Ir—P'	116.3 (2)
Ir—N	1.771 (12)	P—Ir—N	100.4 (3)
Ir—O1	2.954 (11)	P—Ir—N'	93.2 (3)
P—N	3.178 (11)	Ir—N—O1	163.5 (10)
P—N'	3.011 (10)	N—Ir—N'	154.2 (7)
N—N'	3.452 (25)	P—Ir—O1	97.5 (2)
N—O1	1.213 (13)	P—Ir—O1'	89.2 (2)
P—C (av of 3)	1.814 (7)	O1—Ir—O1'	167.5 (5)
P—P'	3.973 (6)	Ir—P—R1C1	114.2 (3)
Ir—H(R2H6)	3.08	Ir—P—R2C1	110.4 (2)
N—H(R1H2)	2.80	Ir—P—R3C1	114.0 (3)
N—H(R2H6)'	2.61	R1C1—P—R2C1	104.7 (3)
O1—H(R1H2)	2.67	R1C1—P—R3C1	108.1 (4)
O1—H(R2H6)'	3.02	R2C1—P—R3C1	104.6 (3)
—Intermolecular dist, Å—		Dihedral angle between	
N—H(R1H4)	3.04	P—Ir—P' and N—Ir—N'	85.7 (3)
O1—H(R3H2)	2.79		

^a X and X' represent atoms related by the twofold axis of symmetry.

TABLE V

ROOT-MEAN-SQUARE AMPLITUDES OF VIBRATION (Å) ^a			
Atom	Min	Intermed	Max
Ir	0.1879 (10)	0.1966 (11)	0.2315 (10)
P	0.186 (4)	0.194 (4)	0.204 (4)
Cl	0.217 (7)	0.287 (6)	0.351 (7)
N	0.208 (14)	0.232 (14)	0.325 (14)
O1	0.225 (14)	0.297 (12)	0.449 (14)

^a Measured along the principal axes of the thermal ellipsoids.

$(\text{SO}_2)(\text{P}(\text{C}_6\text{H}_5)_3)_2$,¹⁶ 2.359 (9) and 2.328 (8) Å, but it is shorter than those found for $\text{IrI}(\text{O}_2)(\text{CO})(\text{P}(\text{C}_6\text{H}_5)_3)_2$,¹⁵ 2.391 (8) and 2.386 (8) Å, and $[\text{IrCl}(\text{CO})(\text{NO})(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{BF}_4]$, 2.408 (3) and 2.407 (3) Å.¹ The Ir—N and N—O distances are 1.771 (12) and 1.213 (13) Å, respectively. The bond angles around iridium are: N—Ir—N', 154.2 (7)°; P—Ir—P', 116.3 (2)°; P—Ir—N', 93.2 (3)°; P—Ir—N, 100.4 (3)°. The dihedral angle between the planes P—Ir—P' and N—Ir—N' is 85.7 (3)°.

As one of the aims of this work was to define the oxidation state of the metal by examining the stereochemistry of the complex, some results obtained by other workers for similar d^{10} and d^8 complexes will be described.

The isoelectronic d^{10} complex $\text{Pt}(\text{CO})(\text{P}(\text{C}_6\text{H}_5)_3)_3$ shows only small distortions from the ideal tetrahedral geometry (P—Pt—P = 104.7 (5)° and P—Pt—C = 113.7 (5)° in a trigonal space group; P—Pt—P = 111.4 (3)° and P—Pt—C = 107 (1)° in a monoclinic space group).¹⁹ However, the complexes $\text{Cu}(\text{NO}_3)(\text{P}(\text{C}_6\text{H}_5)_3)_2$ (P—Cu—P = 131.1°),²⁰ $\text{Cu}(\text{BH}_4)(\text{P}(\text{C}_6\text{H}_5)_3)_2$ (P—Cu—P = 123.3°),²¹ and the phosphine oxide and arsine oxide complexes $[\text{CuCl}(\text{OP}(\text{C}_6\text{H}_5)_2)_2\text{C}_2\text{H}_4]_n$ (Cl—Cu—O = 140.2°)²² and $\text{HgCl}_2(\text{OAs}(\text{C}_6\text{H}_5)_3)_2$ (Cl—Hg—Cl = 146.6°, O—Hg—O = 92.5°)²³ show large distortions from the ideal. In fact many four-coordinate complexes of mercury(II) show large distortions from the ideal tetrahedral geometry.²⁴

In general *trans* square-planar d^8 complexes of the platinum metals show smaller distortions from the ideal. For example the following bond angles have been reported for *trans*- $\text{RhClClCS}(\text{P}(\text{C}_6\text{H}_5)_3)_2$,²⁵ P—Rh—Cl, 89.4 (1) and 87.3 (1)°; P—Rh—C, 91.6 (2) and 92.0 (2)°; Cl—Rh—C, 174.5 (6)°; P—Rh—P, 175.6 (1)°. Similar bond angles have been reported for *trans*- $[\text{PtCl}(\text{CO})(\text{P}(\text{C}_2\text{H}_5)_3)_2][\text{BF}_4]$.²⁶

The bond angles observed for $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$ (see Table IV) are intermediate between those expected for ideal tetrahedral and square-planar geometries. From a consideration of bond angles it is impossible to decide unambiguously on the oxidation state of the metal because d^{10} tetrahedral complexes often show large distortions from the ideal (see above), and although the square-planar d^8 complexes which have been studied show only small distortions, the effect of two electrons in antibonding nitrosyl orbitals is difficult to evaluate.

However, the following factors lead us to favor the d^{10} formulation. The iridium—nitrogen bond length (1.771 (12) Å) is 0.2 Å shorter than that found for $[\text{IrCl}(\text{CO})(\text{NO})(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{BF}_4]$,^{1,2} where the nitrosyl

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ligand is behaving as a Lewis acid, but is in the range found for NO^+ complexes (1.57–1.78 Å; see Table VI of ref 2). The Ir–N–O bond angle ($163.5 (10)^\circ$), though far from 180° , is just within the range observed for NO^+ complexes (see Table VI of ref 2). Finally the large N–Ir–N' bond angle can perhaps be accounted for by nonbonded interactions. The nonbonded contacts between the nitrosyl group and the *ortho*-hydrogen atoms of the triphenylphosphine groups are near those predicted from van der Waals radii.²⁷ We have calculated that these contacts would decrease considerably if the N–Ir–N' bond angle were $109^\circ 28'$. The importance of these steric factors might be evaluated if the crystal structure of $\text{Pt}(\text{CO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2$ were determined.

We have made many attempts to isolate crystals of the five-coordinate complex $\text{IrI}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2$.

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However each time mixtures of $\text{IrI}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2$ and $\text{IrI}_2(\text{NO})(\text{P}(\text{C}_6\text{H}_5)_3)_2$ have been obtained. We wish to speculate that $[\text{Ir}(\text{NO})_2(\text{P}(\text{C}_6\text{H}_5)_3)_2][\text{ClO}_4]$ is able to react with lithium halides to give these unusual five-coordinate complexes for two reasons. First, the geometry of the cation is intermediate between tetrahedral and square planar and therefore the activation energy for formation of a five-coordinate complex is smaller than that for a relatively undistorted tetrahedral complex such as $\text{Pt}(\text{CO})(\text{P}(\text{C}_6\text{H}_5)_3)_3$. Second, nitrosyl is an amphoteric ligand and the two electrons which exceed the inert-gas rule can be accommodated in antibonding nitrosyl orbitals rather than in metal–ligand σ -antibonding orbitals. The fact that nitrosyl is able to behave like an electron well in such reactions suggests some very interesting catalytic possibilities.

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The Crystal and Molecular Structure of Hexakis(urea)titanium(III) Iodide

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The crystal structure of hexakis(urea)titanium(III) iodide, $\text{Ti}[\text{OC}(\text{NH}_2)_2]_6\text{I}_3$, has been determined from three-dimensional X-ray data collected by counter methods. The structure has been refined to a final unweighted R factor of 0.067 and a weighted R factor of 0.064 for 499 independent nonzero reflections. The salt crystallizes in the rhombohedral space group $R\bar{3}c$ with six formula units in a hexagonal unit cell of dimensions $a = 17.67 (2) \text{ \AA}$ and $c = 14.15 (2) \text{ \AA}$. The structure consists of octahedral complex ions having rigorous D_3 point symmetry and essentially linear arrays of iodide ions parallel to the crystallographic threefold axis. The titanium–oxygen bond distance is $2.014 (5) \text{ \AA}$. There are significant distortions of the coordination polyhedron from the idealized O_h symmetry and from the D_{3d} model often assumed for trigonally distorted octahedral complexes. The three independent O–Ti–O angles are $90.5 (2)$, $86.4 (3)$, and $92.8 (3)^\circ$, and the angles between successive pairs of Ti–O bonds viewed in projection along the threefold axis are 54.4 and 65.6° . The stability of this complex to aerial oxidation apparently arises from intramolecular hydrogen bonding involving the urea molecules and also hydrogen bonds involving the iodide ions. The d-orbital energy level sequence for a crystal (ligand) field of D_3 symmetry, where the inclination of metal–ligand bond to the threefold axis is 54.74° , can be described entirely in terms of the parameter Dq , and the ground state for a d^1 complex ion is 2A_1 .

Introduction

Coordination compounds of titanium(III) have long occupied a prominent place in the development at various levels of sophistication of crystal and ligand field theory. Despite this fact, however, there is a paucity of accurate structural information from X-ray measurements on their detailed stereochemistry in the solid, principally because the majority of such compounds are highly unstable with respect to oxidation and many hydrolyze rapidly in the presence of moisture. The isolation and handling of suitable single crystals is thus a difficult operation.

One of the more stable species is the hexakis-urea complex ion which was first isolated as the perchlorate

salt² and subsequently as the iodide.³ One of the interesting features of the latter compound is its marked stability relative to other titanium(III) species, and a preliminary report of the crystal structure, in the form of a two-dimensional projection, has already been given.⁴ This preliminary work did not reveal the origin of the stability of the complex, however, and we have therefore carried out a more detailed study, both to investigate this point and to assess the nature of the distortion from regular octahedral symmetry, which is indicated by the observed variation

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