

***BMO* ON STRONGLY PSEUDOCONVEX DOMAINS: HANKEL OPERATORS, DUALITY AND $\bar{\partial}$ -ESTIMATES**

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ABSTRACT. We study the condition that characterizes the symbols of bounded Hankel operators on the Bergman space of a strongly pseudoconvex domain and show that it is equivalent to *BMO* plus analytic. (Here we mean the Bergman metric *BMO* of Berger, Coburn and Zhu.) In the course of the proof we obtain new $\bar{\partial}$ -estimates that may be of independent interest. Some applications include a decomposition of *BMO* similar to the classical $L^\infty + \widetilde{L^\infty}$, and two characterizations of the dual of *VMO* (which is also a predual of *BMO*).

In addition, we obtain some partial results on the boundedness of Hankel operators in L^1 norm.

1. INTRODUCTION

Let D be a strongly pseudoconvex domain in \mathbb{C}^n , and let dV denote $2n$ -dimensional volume measure on D . (The reader may assume on first reading that D is \mathbb{D} , the open unit disk, substituting area measure dA for dV .) Let $L^2 = L^2(dV)$ denote the space of square integrable measurable functions on D , \mathcal{H} the space of all holomorphic functions on D , and $A^2 = L^2 \cap \mathcal{H}$. Denote by P the orthogonal projection from L^2 onto A^2 , which is effected by integrating against the Bergman kernel $K = K_D(z, w)$. Define the *Hankel operator* H_f with symbol f by

$$H_f g = fg - P(fg)$$

where f is some L^1 function on D . This is defined (via the integral formula for P) whenever g is a bounded analytic function on D . Moreover, if $h \in A^1 = L^1 \cap \mathcal{H}$, then $H_h g = 0$ for all bounded analytic g . This allows H_f to be densely defined for any $f \in L^1 + \mathcal{H}$, and it is guaranteed to be defined on all of A^2 (and bounded) whenever f is bounded. It may, however, extend to be a bounded linear operator on A^2 (or even A^p , $1 < p < \infty$) even when f is not bounded.

Necessary and sufficient conditions for the extension of H_f to a bounded operator from A^p to L^p were obtained in [Lu3] for the unit disk, and these results were extended to strongly pseudoconvex domains and other settings in [Li2]. The conditions are formally the same in the two settings and will be discussed in more detail below. For now, let us say that the condition for H_f to be bounded from A^p to L^p is a uniform bound on the distance from $f|_{B(z)}$

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to analytic functions, where $B(z)$ is a ball in the Bergman metric (of some fixed size) centered at z , and the distance is taken in the L^p norm with respect to normalized Lebesgue measure on $B(z)$. We informally termed this property the *BDA* condition (for “Bounded Distance to Analytic”) but did not formally define a space of functions based on it, perhaps because we were convinced it had a better characterization.

Part of the purpose of this paper is to present such a better characterization. If we note that the space *BMO* of Berger, Coburn, and Zhu [BCZ] is precisely the space with “bounded distance to constants”, it is clear that $BMO + \mathcal{H}$ is contained in *BDA*. Our main purpose here is to prove the opposite containment and to explore some of the implications of this fact and of its proof. These implications include a representation of *BMO* similar to the “ $L^\infty + \widetilde{L^\infty}$ ” representation of the classical *BMO* on \mathbb{R}^n and several representations of the predual of *BMO* (which is also the dual of *VMO*). In addition, some partial results on the boundedness of Hankel operators from A^1 to L^1 will be presented, including a complete characterization for conjugate analytic symbol.

2. BACKGROUND AND HISTORY

The seminal paper on the boundedness of H_f is S. Axler’s [Ax2]. In that paper f is assumed to be the conjugate of an analytic function h in the unit disk \mathbb{D} , $f = \bar{h}$. The result obtained is that $H_{\bar{h}}$ extends to be a bounded operator from A^2 to L^2 if and only if h is in the *Bloch space* \mathcal{B} , that is, $(1 - |z|^2)h'(z)$ is bounded in \mathbb{D} . Moreover, H_f is compact if and only if $h \in \mathcal{B}_0$, that is, $(1 - |z|^2)h'(z) \rightarrow 0$ as $|z| \rightarrow 1$. Other results (also in [Ax2]) show that \mathcal{B} is $\mathcal{H} \cap BMO$ and \mathcal{B}_0 is $\mathcal{H} \cap VMO$.

The next significant steps were taken by D. Békollé, C. Berger, L. Coburn and K. Zhu in the two papers [BCZ] and [BBCZ]. In those papers the setting is a bounded symmetric domain in \mathbb{C}^n . They showed the following:

Theorem A (Békollé, Berger, Coburn, Zhu). *The following are equivalent for a bounded symmetric domain.*

- (a) H_f and $H_{\bar{f}}$ are simultaneously bounded.
- (b) f belongs to *BMO*.

Let us interrupt here to define *BMO* and the related *VMO*. First, let β denote the Bergman distance (the integrated form of the Bergman metric) on the bounded domain D . The symbol D may now represent either a bounded symmetric domain or a strongly pseudoconvex domain. On the disk \mathbb{D} , β is the same as the hyperbolic distance. Let $r > 0$ be some fixed number and let $B(z) = B(z, r) = \{w \in D : \beta(z, w) < r\}$ denote the ball with respect to β centered at z with radius r . Finally, let the Lebesgue measure of a set E be denoted by $|E|$.

2.1. Definition. Let $1 < p < \infty$ and let f be a measurable function on D locally in L^p . The *p-mean oscillation* of f at z is defined to be

$$MO_p(f, z) = \left(\inf_{c \in \mathbb{C}} \frac{1}{|B(z)|} \int_{B(z)} |f - c|^p dV \right)^{1/p}.$$

We say that $f \in BMO_p$ (respectively, $f \in VMO_p$) if and only if

$$\sup_{z \in D} MO_p(f, z) < \infty \quad \left(\text{respectively, } \lim_{z \rightarrow \partial D} MO_p(f, z) = 0 \right)$$

and we write $\|f\|_{BMO, p}$ for the supremum. If the subscript p is omitted in any of the above notation then $p = 2$ is assumed.

It is well known that the different choices for r yield equivalent norms on BMO_p . Theorem A above is accompanied by a similar one in which “compact” replaces “bounded” and “ VMO ” replaces “ BMO ”.

Theorem A is equivalent to characterizing the bounded Hankel operators for real f only. It also extends Axler’s result, since $H_{\bar{f}} = 0$ when \bar{f} is analytic and $\mathcal{B} = BMO \cap \mathcal{H}$ ([Ax2]). Nevertheless, it does not yield a characterization of the boundedness of H_f alone without additional restrictions on f .

Theorem A was extended to strongly pseudoconvex domains by the first author (H. Li) in [Li1] and independently by F. Beatrous and Song-Ying Li (in [BL]). Then in [Lu3] the second author (Luecking) was able to remove all restrictions on f (at least for the unit disk), requiring only that f be integrable. The resulting characterization utilized an approach that did not quite work in other settings. The difficulties were overcome in [Li2] by the first author. That paper extended the result to Hankel operators on the Bergman space of a strongly pseudoconvex domain in \mathbb{C}^n , but also extended the proof to all the usual weighted Bergman spaces and to the Fock space. A different method (from [Lu3]) was used for part of the proof, and this improved the result even in the unit disk. Unfortunately, the approach has not yet yielded a proof for bounded symmetric domains, so that Theorem A is still the best result known in that setting.

Before proceeding further, we should formally define a few of the terms we have already used.

2.2. Definition. A domain D in \mathbb{C}^n will be called *strongly pseudoconvex* if there is a strictly plurisubharmonic defining function ρ for D . That is, ρ is defined and smooth (i.e., C^∞) in a neighborhood of \bar{D} , the Levi form for ρ is strictly positive definite on \bar{D} , the differential $d\rho$ is nonvanishing on ∂D , and $D = \{z : \rho(z) < 0\}$. The *Bloch space* \mathcal{B} (respectively, the *little Bloch space* \mathcal{B}_0) of the strongly pseudoconvex domain D is the collection of all functions f holomorphic on D satisfying

$$(1) \quad \sup_{z \in D} |\rho(z) \partial f(z)| < \infty \quad \left(\text{respectively, } \lim_{z \rightarrow \partial D} \rho(z) \partial f(z) = 0 \right).$$

The interpretation of “smooth” to mean “ C^∞ ” can be replaced with the interpretation “as differentiable as necessary”. We will not be concerned with the exact level of differentiability required.

It automatically follows from (1) that the components of ∂f in the complex directions perpendicular to $\partial \rho$ are $O(|\rho|^{-1/2})$ (see [KM]). Moreover, it remains true in strongly pseudoconvex domains, as in the unit disk, that $\mathcal{B} = \mathcal{H} \cap BMO_p$ and $\mathcal{B}_0 = \mathcal{H} \cap VMO_p$ for all $p \geq 1$ (see [Li1]).

Let us turn to the characterizations obtained in [Lu3] and [Li2]. Recall that $r > 0$ has been fixed and $B(z) = B(z, r)$. Fix also $p > 1$ and measurable f on D , and define the following local measurement of the distance to analytic

functions:

$$(2) \quad F(z) = \left(\inf_{h \in \mathcal{H}} \frac{1}{|B(z)|} \int_{B(z)} |f - h|^p dV \right)^{1/p}.$$

The function $F(z)$ depends on f and on p , but this will not be reflected in the notation. Of course, F also depends on r , but this fact is inessential.

Theorem B. *Let D be a strongly pseudoconvex domain in \mathbb{C}^n with strictly plurisubharmonic defining function ρ (satisfying $d\rho \neq 0$ on ∂D), and let $f \in L^1(D)$. For $p > 1$ the following are equivalent.*

- (a) H_f extends to a bounded operator from A^p to L^p .
- (b) F is bounded on D for some (respectively, every) choice of r .
- (c) $f = f_1 + f_2$, where

$$G(z) \stackrel{\text{def}}{=} \left(\frac{1}{|B(z)|} \int_{B(z)} |f_1|^p dV \right)^{1/p}$$

is bounded on D for some (respectively, every) choice of $r > 0$ and f_2 is a C^∞ function such that

$$(3) \quad |\rho(z)| |\bar{\partial} f_2(z)| + |\rho(z)|^{1/2} |\bar{\partial} f_2(z) \wedge \bar{\partial} \rho(z)|$$

is also bounded on D .

In the unit disk $\rho(z) = |z|^2 - 1$, the first term in (3) is $(1 - |z|^2) \partial f_2 / \partial \bar{z}$, and the second term is always 0. Thus $\bar{\partial} f_2$ grows at most inversely as the distance to the boundary. This is also the case in a strongly pseudoconvex domain, but there the second term imposes the additional condition that the component of $\bar{\partial} f_2$ “parallel to the boundary” grows inversely as the square root of the distance to the boundary. The combination of derivatives in (3) appears in [V, Part 2] in connection with the study of BMO (on the boundary) and the $\bar{\partial}$ -equation.

Condition (b) of the theorem will be referred to as the “ BDA condition”, and it will be convenient to define now the space BDA_p :

2.3. Definition. Let f be a measurable function on D and define F as in (2). If F is bounded, we say that f is in BDA_p .

The condition $f \in BDA_p$ clearly implies that f is locally in L^p . Moreover, the implication (b) \Rightarrow (c) of Theorem B does not actually require the prior assumption that $f \in L^1$, and the proof of the implication (c) \Rightarrow (a) shows that if $f \in BDA_p$ then $f \in L^p + \mathcal{H}$. There is an obvious seminorm $\|f\|_{BDA_p} = \sup_{z \in D} F(z)$ on BDA_p which makes BDA_p / \mathcal{H} into a Banach space. The following is our improvement of the BDA condition.

2.4. Theorem. For $p \geq 1$, $BDA_p = BMO_p + \mathcal{H}$.

We will prove this in §3 (delaying the proof of the main lemma until §6). In §4 we will explore some of the implications of this theorem and of its proof on the predual of BMO_p . In §5 we will consider some partial results on the boundedness of Hankel operators in the L^1 norm.

3. THE MAIN THEOREM

Let us begin by recalling a brief sketch of the proof of (c) \Rightarrow (a) in Theorem B. Observe that for $f \in C^1$ and g analytic, $H_f g$ is the minimal L^2 norm

solution of the $\bar{\partial}$ -equation

$$(4) \quad \bar{\partial}u = g\bar{\partial}f.$$

That is, $H_f g = (I - P)u$, where u is any solution of (4). A class of kernels that solved the $\bar{\partial}$ -equation in a strongly pseudoconvex domain was described in [BA]. A particular one of these $L(z, w)$ was constructed in [Li2]. $L(z, w)$ is an $(n, n-1)$ -form in w such that $u(z) = \int \varphi(z) \wedge L(z, w)$ solves $\bar{\partial}u = \varphi$ for any $\bar{\partial}$ -closed $(0, 1)$ -form φ which satisfies

$$(5) \quad |\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho| \in L^1.$$

Moreover, it was shown in [Li2] that if the expression in (5) is bounded then the kernel $\varphi(w) \wedge L(z, w)$ defines a bounded operator on L^p , $p > 1$. Since also $I - P$ is bounded on L^p , this result and condition (c) of Theorem B imply that $H_{f_2}g = (I - P) \int g(w) \bar{\partial}f_2(w) \wedge L(\cdot, w)$ is bounded from A^p to L^p .

In order to prove Theorem 2.4 we will construct another such kernel L and show that the same condition (the boundedness of (5) above) implies that $\varphi(w) \wedge L(z, w)$ defines a bounded map from L^∞ to BMO_p . Having accomplished this we can finish the proof in short order as follows: If $f \in BDA_p$ then by the (b) \Rightarrow (c) portion of Theorem B there is a decomposition $f = f_1 + f_2$. The condition (c) on f_1 clearly places it in BMO_p . Applying the operator with kernel $\bar{\partial}f_2(w) \wedge L(z, w)$ to the function 1 gives an element u of BMO_p that solves $\bar{\partial}u = \bar{\partial}f_2$. Then $f_2 - u \in \mathcal{H}$ and so $f = (f_1 + u) + (f_2 - u) \in BMO_p + \mathcal{H}$. We will actually show somewhat more. The solution u will be shown to have bounded oscillation rather than just bounded mean oscillation, and, by showing that P is bounded from BMO_p to \mathcal{B} , we will see that if $f \in L^1$ and H_f is bounded on A^p then $(I - P)f \in BMO_p$. (One minor point we ought to note is that the $p = 1$ version of this argument is valid: As will be remarked in Theorem 5.1, the forward implications (a) \Rightarrow (b) \Rightarrow (c) of Theorem B follow in the same way for $p = 1$ as for $p > 1$, and Proposition 3.2 will show that P is bounded also from BMO_1 to \mathcal{B} .)

In the case of the unit disk, the minimal L^2 norm solution to the $\bar{\partial}$ -equation $\bar{\partial}u = \varphi d\bar{z}$ is

$$u(z) = \frac{1}{2\pi i} \int_{\mathbb{D}} \frac{1 - |w|^2}{(z - w)(1 - \bar{z}\bar{w})} \varphi(w) d\bar{w} \wedge dw.$$

Following the above approach, $H_{f_2}g$ is given by

$$H_{f_2}g(z) = \frac{1}{2\pi i} \int_{\mathbb{D}} \frac{g(w)(1 - |w|^2)\bar{\partial}f_2(w)}{(z - w)(1 - \bar{z}\bar{w})} \wedge dw.$$

Letting $T(z, w) = [2\pi(z - w)(1 - \bar{z}\bar{w})]^{-1}$, it is straightforward to prove that T defines a bounded operator on the L^p spaces. This is the (nearly trivial) unit disk version of the work in [Li2]. Given the explicit nature of the kernel T , it is another straightforward exercise to show that it takes bounded functions to BMO_p . In fact it is easy to show more, namely that T takes bounded functions to functions u that satisfy

$$|u(z_1) - u(z_2)| \leq C(1 + \beta(z_1, z_2))$$

with C independent of z_1 and z_2 . We will obtain a similar result for strongly pseudoconvex domains.

3.1. **Lemma.** *Let φ be a $\bar{\partial}$ -closed $(0, 1)$ -form satisfying*

$$|\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho| \text{ is bounded.}$$

Then there is a solution u to the equation $\bar{\partial}u = \varphi$ which has the form $u = u_1 + u_2$, where

$$\|u_1\|_\infty + \|\rho\nabla u_2\|_\infty \leq C\| |\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho| \|_\infty$$

and

$$|\partial u_2(z) \wedge \partial \rho(z)| + |\bar{\partial} u_2(z) \wedge \bar{\partial} \rho(z)| \leq C|\rho(z)|^{-1/2}\| |\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho| \|_\infty.$$

Note that the gradient of u_2 grows as the reciprocal of the distance to the boundary but has components in the complex directions parallel to the boundary which are of size $C|\rho|^{-1/2}$. This implies that $|u(z_1) - u(z_2)| \leq C$ whenever $z_2 \in B(z_1)$, which clearly implies that $u \in BMO_p$ for any $p \geq 1$.

The proof of Lemma 3.1 is postponed until §6.

As described earlier, this lemma is sufficient to prove Theorem 2.4. For later use, remember that what will actually be proven is that $u = u_1 + u_2$ where $u_1 \in L^\infty$ and the gradient of u_2 behaves like the derivative of a Bloch function. In particular $f \in BDA_p$ implies that $f = f_1 + g + h$, where f_1 satisfies condition (c) of Theorem B, h is analytic, and $|g(z_1) - g(z_2)| \leq C\beta(z_1, z_2)$ for all z_1 and z_2 in D . If $f \in BMO_p$ then the same is true with $h \equiv 0$ (because then $h \in \mathcal{B}$ and it can be absorbed in the g).

We now show that the Bergman projection P is a bounded operator from BMO_p to \mathcal{B} that moreover takes VMO_p to \mathcal{B}_0 .

3.2. **Proposition.** *Let $p \geq 1$. If $f \in BMO_p$ then $Pf \in \mathcal{B}$ (the Bloch space) and so $(I-P)f \in BMO_p$. If $f \in VMO_p$ then $Pf \in \mathcal{B}_0$ and $(I-P)f \in VMO_p$.*

Proof. If $f \in BMO_p$ then $f = f_1 + g$ where f_1 satisfies condition (c) in Theorem B, that is,

$$\sup_{z \in D} \left(\frac{1}{|B(z)|} \int_{B(z)} |f_1|^p dV \right)^{1/p} < \infty,$$

while g satisfies a Lipschitz condition:

$$|g(z_1) - g(z_2)| \leq C\beta(z_1, z_2).$$

We first show that $Pg \in \mathcal{B}$. By Definition 2.2, for this it suffices to show that $\rho(z)\partial(Pg)(z)$ is bounded. Using Fefferman's formula [Fe] for the Bergman kernel $K(z, w)$, it follows that

$$|\partial K(z, w)/\partial z_k| \leq C|K(z, w)||\Psi(z, w)|^{-1} \leq C\Psi(z, w)^{-n-2},$$

where $\Psi(z, w)$ is defined in §5. (It is essentially the same as $X(z, w)$ in [Fe] and $\widehat{\Phi}(z, w)$ in [R].) Now, using estimates given in [Li1, Theorem 4 and Corollary 11], we obtain $\beta(z, w) \leq C_\varepsilon|\rho(z)|^{1-\varepsilon}|\rho(w)|^{-\varepsilon}|K(z, w)|^{-2\varepsilon/(n+1)}$ for

sufficiently small ε . Thus

$$\begin{aligned} \left| \frac{\partial(Pg)}{\partial z_k} \right| &= \left| \int_D g(w) \frac{\partial K(z, w)}{\partial z_k} dV(w) \right| \\ &= \left| \int_D (g(w) - g(z)) \frac{\partial K(z, w)}{\partial z_k} dV(w) \right| \\ &\leq C \int_D \beta(z, w) |K(z, w)| |\Psi(z, w)|^{-1} dV(w) \\ &\leq C_\varepsilon \int_D |\rho(z)\rho(w)|^{-\varepsilon} |\Psi(z, w)|^{-n-2+2\varepsilon} dV(w) \\ &\leq CC_\varepsilon / |\rho(z)|, \end{aligned}$$

where the last inequality comes from [R, Lemma 7.3, p. 310]. This is what was required for g .

It is clear from the condition on f_1 and Hölder's inequality that the function

$$\frac{1}{|B(z)|} \int_{B(z)} |f_1| dV$$

is bounded. For each $z \in D$ let $E(z)$ be a polydisk centered at z with radii $c|\rho(z)|$ in the complex directions normal to the boundary and $c\sqrt{|\rho(z)|}$ in the complex directions parallel to the boundary. Choose c so that $E(z) \subset B(z)$ for all z and recall that $\inf_{w \in B(z)} |E(w)| \sim |B(z)|$. Now $K(z, w) = (1/|E(w)|) \int_{E(w)} K(z, \zeta) dV(\zeta)$, so that by Fubini's theorem

$$\begin{aligned} Pf_1(z) &= \int_D f_1(w) K(z, w) dV(w) \\ &= \int_D \left[\int_D f_1(w) \frac{\chi_{E(w)}(\zeta)}{|E(w)|} dV(w) \right] K(z, \zeta) dV(\zeta). \end{aligned}$$

Since $\chi_{E(w)}(\zeta) \leq \chi_{B(w)}(\zeta) = \chi_{B(\zeta)}(w)$, it is clear that the expression in brackets above is a bounded function. Thus Pf_1 , being the projection of a bounded function, belongs to \mathcal{B} . (See [C].)

The VMO_p case is entirely similar. \square

Combining Theorem B, Theorem 2.4 and Proposition 3.2 gives us the following:

3.3. Theorem. *Let $f \in L^1$ and $p > 1$. Then H_f is bounded from A^p to L^p if and only if $(I - P)f \in BMO_p$.*

Proof. If $(I - P)f \in BMO_p$ then $Pf \in L^1$ and $H_{Pf} = 0$. Thus $H_f = H_{(I-P)f}$ is bounded. Conversely, if H_f is bounded then, by Theorem B, f belongs to BDA_p and, by Theorem 2.4, $f \in BMO_p + \mathcal{H}$. The holomorphic part will lie in L^1 , and $I - P$ applied to it will give 0. The conclusion now follows by Proposition 3.2. \square

4. BMO AND RELATED SPACES

Perhaps as interesting as the main theorem are some of its consequences. In contrast to the classical VMO on \mathbb{R} or $\partial\mathbb{D}$, the balls $B(z)$ over which the oscillation is averaged do not vary in size and it is only at the boundary ("at

infinity" so to speak) that the mean oscillation conditions are of consequence. Moreover, the best that can be said of the local behavior of a function in BMO_p or VMO_p is that it is locally in L^p . One consequence of this difference is the fact that the spaces VMO_p and BMO_p vary with p . Another consequence is that the analogy between BMO and L^∞ in the classical setting is replaced by an analogy between BMO_p and L_p^∞ , where L_p^∞ will be defined below. Moreover, the continuous functions no longer play the same rôle in the theory. Nevertheless the fact that H_f is bounded on L^p if and only if $f \in BMO_p + \mathcal{H}$ and compact if and only if $f \in VMO_p + \mathcal{H}$ suggests that BMO and VMO play similar rôles in both settings. This is reinforced by the decomposition and duality results to follow.

4.1. Definition. Let $0 < p, q \leq \infty$ and let f be a measurable function on \mathbb{D} . We say that f belongs to L_q^p if the function

$$z \mapsto \left(\frac{1}{|B(z)|} \int_{B(z)} |f|^q dV \right)^{1/q}$$

belongs to L^p . If $q = \infty$ we naturally replace the above function with

$$z \mapsto \operatorname{ess\,sup}_{w \in B(z)} |f(w)|.$$

We give L_q^p the obvious "norm", and it is a Banach space when p and q are both at least 1. In case $1 \leq p, q < \infty$ the dual of L_q^p is $L_{q'}^{p'}$ with the pairing $\langle f, g \rangle = \int_{\mathbb{D}} f \bar{g} dV$. In this notation, the condition on f_1 in part (c) of Theorem B is that $f_1 \in L_p^\infty$.

The following result will be stated in the context of the unit disk \mathbb{D} in \mathbb{C} . Afterwards, corresponding results in a strongly pseudoconvex domain in \mathbb{C}^n will be discussed briefly.

Recall that T is the integral operator on L^p of \mathbb{D} whose kernel is $T(z, w) = [\pi(z-w)(1-\bar{w}z)]^{-1}$, and P is the Bergman projection whose kernel is $K(z, w) = [\pi(1-\bar{w}z)]^{-2}$.

4.2. Theorem. If $f \in BMO_p$ then there exist functions $f_1 \in L_p^\infty$, $\varphi \in L^\infty$ and $\psi \in L^\infty$ such that

$$(6) \quad f = f_1 + T\varphi + P\psi.$$

The functions can be found to satisfy estimates

$$(7) \quad (1/C)\|f\|_{BMO,p} \leq \|f_1\|_{L_p^\infty} + \|\varphi\|_\infty + \|\psi\|_\infty \leq C\|f\|_{BMO,p}.$$

Proof. Let $f \in BMO_p$. Clearly $f \in BDA_p$, and so $f = f_1 + f_2$ as in part (c) of Theorem B. We have seen that $(I-P)f_2 = T(\varphi)$ with $\varphi(w) = \bar{\partial} f_2(w)(1-|w|^2)$, so $f = f_1 + T\varphi + h$ with h holomorphic and φ bounded. Clearly h is also in BMO_p , and so the results in [Ax2] show that $h \in \mathcal{B}$, whence (see [Ax1]) $h = P\psi$ with ψ bounded as required. The construction in [Lu3] gives an estimate on the first and second term of the sum in (7). The results in [Ax1] give an estimate on the third. Combined, these give the right-hand inequality of (7). Given a function f of the form (6), the boundedness of T and P from L^∞ to BMO_p gives the left-hand inequality. \square

In a strongly pseudoconvex domain the following formulation seems the best we can do for now.

4.3. Theorem. If $f \in BMO_p$ then there exists functions f_1, f_2 , and ψ such that $f_1 \in L_p^\infty$, f_2 satisfies the conditions of part (c) of Theorem B, ψ is bounded, and

$$(8) \quad f = f_1 + \int_D \bar{\partial} f_2(w) \wedge L(\cdot, w) + P\psi,$$

where L is the kernel for the solution of the $\bar{\partial}$ -equation constructed in §6. The functions can be found to satisfy estimates

$$(9) \quad (1/C)\|f\|_{BMO,p} \leq \|f_1\|_{L_p^\infty} + \|\rho \bar{\partial} f_2\| + |\rho|^{1/2} \|\bar{\partial} f_2 \wedge \bar{\partial} \rho\|_\infty + \|\psi\|_\infty \\ \leq C\|f\|_{BMO,p}.$$

Proof. The proof goes just like the proof of Theorem 4.2 except that Coupet's work [C] is cited instead of Axler's and Li's [Li2] instead of Luecking's \square

In the statement of Theorem 4.3 any kernel that solves the $\bar{\partial}$ -equation and satisfies the estimates of Lemma 6.4 may replace L .

We now define a version, suitable for our setting, of atomic H^1 . In a sense, this goes to the opposite extreme from both the standard atomic theory and the "local" atomic theory of D. Goldberg ([Go]) in that the support of our atoms is not allowed to vary in size.

4.4. Definition. Let $q > 1$. A q -atom is a measurable function $a(z)$, locally in L_q , such that

(a) a is supported by some ball $B(z_a)$ (r is fixed),

(b) $\int_D a dV = 0$,

(c) $((1/|B(z_a)|) \int_{B(z_a)} |a|^q dV)^{1/q} = |B(z_a)|^{-1}$.

The space H_q^1 is defined to consist of all measurable functions f which can be represented as an l^1 sum of atoms: $f = \sum_j \lambda_j a_j$ with $\sum_j |\lambda_j| < \infty$. The norm of a function in H_q^1 is defined to be the infimum of the sums $\sum |\lambda_j|$ over all such representations of f .

It is easy to verify that there is a constant $C = C(q, r)$ such that a q -atom has L_q^1 norm greater than $1/C$ and less than C so that $H_q^1 \subset L_q^1$ and $\|f\|_{L_q^1} \leq C\|f\|_{H_q^1}$.

4.5. Proposition. Let $1 < p, q < \infty$ with $\frac{1}{p} + \frac{1}{q} = 1$. Then $VMOP_p^* = H_q^1$ and $(H_q^1)^* = BMO_p$.

The pairing is the sesquilinear one $\langle f, g \rangle = \int_D f \bar{g} dV$ and is to be taken in the following sense: Functions with compact support in D are dense in H_q^1 and the integral makes sense when f is such a function. For arbitrary functions in H_q^1 the integral is obtained as a limit of the compactly supported case.

Proof. Let $g = \sum \lambda_j a_j \in H_q^1$, where each a_j is a q -atom supported in a ball B_j and chosen so that $\sum |\lambda_j| \leq (1 + \varepsilon)\|g\|_{H_q^1}$. For a function $f \in VMOP_p$ choose constants c_j so that

$$\left(\frac{1}{|B_j|} \int_{B_j} |f - c_j|^p dV \right)^{1/p} \leq (1 + \varepsilon)\|f\|_{BMO,p}.$$

Then we have

$$\begin{aligned}
 |\langle f, g \rangle| &= \left| \sum_j \lambda_j \int f \bar{a}_j dV \right| = \left| \sum_j \lambda_j \int_{B_j} (f - c_j) \bar{a}_j dV \right| \\
 &\leq \sum_j |\lambda_j| \left(\int_{B_j} |f - c_j|^p dV \right)^{1/p} \left(\int_{B_j} |a_j|^q dV \right)^{1/q} \\
 &\leq \sum_j |\lambda_j| (1 + \varepsilon) \|f\|_{BMO, p} \leq (1 + \varepsilon)^2 \|g\|_{H_q^1} \|f\|_{BMO, p}.
 \end{aligned}$$

Thus g defines a continuous linear functional on VMO_p whose norm is dominated by its H_q^1 norm. If g happens to be a q -atom, let $\|g\|$ denote its H_q^1 norm and let $\|g\|_*$ denote its norm as a linear functional. Then we have $\|g\|_* \leq \|g\| \leq C$ for $C = C(q, r)$. On the other hand $\|g\|_* = \sup |\int_B f \bar{g} dV|$, where B is the ball supporting g in the definition of atom and the supremum is taken over $f \in VMO_p$ with norm 1. Select any function ψ in $L^p(B, dV/|B|)$ with distance to the constants equal to 1. Extend ψ to a function f in VMO_p by defining it to be constant outside B using the constant c which achieves the distance: $\int_B |\psi - c|^p dV/|B| = 1$. Then f is a norm 1 element of VMO_p and the supremum of $\int f \bar{g} dV$ over all such f is $|B|(\int_B |g|^q dV/|B|)^{1/q} = 1$. Thus $1 \leq \|g\|_* \leq \|g\| \leq C$. Moreover, the definitions show that the supremum over all q -atoms g of $|\int f \bar{g} dV|$ is equal to $\|f\|_{BMO, p}$. Thus the q -atoms form a bounded norming subset of the dual of VMO_p , and now standard functional analysis arguments (see [Bo]) show that every element of the dual is an l^1 sum of q -atoms with the appropriate equivalence of norms.

Since different choices for the radius r of $B(z)$ in the definition of VMO_p give the same space with equivalent norms, we see that the same is true of H_q^1 . Thus any function f in $L^1(D, dV)$ which is supported on a ball $B(z, r')$ of any size and satisfies $\int f dV = 0$ will belong to H_q^1 .

The argument at the beginning of the proof also works to show that BMO_p embeds in the dual of H_q^1 with the dual norm dominated by the BMO_p norm. If now Λ is a linear functional on H_q^1 , then for each ball $B = B(z, r)$, Λ is defined on the subset of $L^q(B, dV/|B|)$ of functions of mean 0. The functions in this space having norm $1/|B|$ are q -atoms. Thus Λ is associated to a family of functions $\{g_B\}$ by $\Lambda(f) = \int_B f \bar{g}_B dV/|B|$, each g_B lying in $L^p(B, dV/|B|)$ and having distance from the constants at most $\|\Lambda\|$. If two such balls overlap, it is easy to show that the corresponding functions differ by a constant on the overlap. (Their difference annihilates all functions with mean 0.) The problem now is to construct a globally defined function g whose value on any given ball $B = B(z, r)$ differs from the corresponding g_B by only a constant. For then

$$\begin{aligned}
 \|g\|_{BMO, p}^p &= \sup_z \inf_c \frac{1}{|B(z, r)|} \int_{B(z, r)} |g - c|^p dV \\
 &= \sup_z \inf_c \frac{1}{|B(z, r)|} \int_{B(z, r)} |g_B - c|^p dV \leq \|\Lambda\|^p.
 \end{aligned}$$

To this end fix a point z_0 in D . Let $r < r_1 < r_2 < \dots$ with $r_k \rightarrow \infty$. Then inductively obtain g_1, g_2, \dots such that for every $k \geq 1$, g_k represents Λ on

$B(z_0, r_k)$. Adjust each g_k by a constant if necessary so that $g_k = g_{k-1}$ on $B(z_0, r_{k-1})$. Then g is well defined by $g(z) = g_k(z)$ on $B(z_0, r_k)$ and agrees up to a constant with g_B on every $B = B(z, r)$. Thus, for every q -atom f , $\Lambda(f) = \int f \bar{g} dV$. \square

We now define a space V_q^∞ related to L_q^∞ in the same way that VMO is related to BMO .

4.6. Definition. If $f \in L_q^\infty$ we say that $f \in V_q^\infty$ if $((1/|B(z)|) \int_{B(z)} |f|^q dV)^{1/q} \rightarrow 0$ as $z \rightarrow \partial D$.

Soon we will need the following small lemma, which is to VMO_p as Theorem 4.2 is to BMO_p .

4.7. Lemma. Every $f \in VMO_p$ can be written

$$f = f_1 + T\varphi + P\psi$$

where $f_1 \in V_p^\infty$, $\varphi \in C_0$, and $\psi \in C_0$. The estimates of Theorem 4.2 (inequalities (7)) hold for f , f_1 , φ , and ψ .

Proof. The proof is just like the proof of Theorem 4.2 but is based on the corresponding compactness criteria in [Lu3] instead of the boundedness criteria. Naturally, there is an analogous result for strongly pseudoconvex domains along the lines of Theorem 4.3. \square

Keeping the convention that $1/p + 1/q = 1$, we turn to a closer examination of the space H_q^1 . It arises as the predual of BMO_p , which itself arises as the set of symbols f for which H_f is bounded. Or rather $BMO_p + \mathcal{H}$ does. The estimates of Theorem B and Lemma 3.1 indicate that for a given function f ,

$$\|f\|_{BDA,p} \leq \inf_{h \in \mathcal{H}} \|f - h\|_{BMO,p} \leq C \|f\|_{BDA,p}.$$

Since $BMO_p \cap \mathcal{H} = \mathcal{B}$ (respectively, $VMO_p \cap \mathcal{H} = \mathcal{B}_0$), the BDA_p norm of a function in BMO_p (respectively, VMO_p) coincides with the norm in BMO_p/\mathcal{B} (respectively, VMO_p/\mathcal{B}_0). The predual of BMO_p/\mathcal{B} (respectively, the dual of VMO_p/\mathcal{B}_0) is the subset of functions in H_q^1 that annihilate analytic functions, that is, $\mathcal{B}_0^\perp \stackrel{\text{def}}{=} \{f \in H_q^1 : \int f \bar{g} dV = 0 \text{ for all } g \in \mathcal{B}_0\}$. Now the BMO_p/\mathcal{B} norm of a function f is equivalent to the norm of the operator H_f from A^p to L^p , and this in turn is equivalent to

$$(10) \quad \sup \langle H_f g_1, g_2 \rangle = \sup \int f g_1 \bar{g}_2 dV$$

where g_1 varies through the unit ball of A^p and g_2 varies through the unit ball of $(A^p)^\perp \subset L^q = L_q^q$. Since $A^p \subset L_\infty^p$ it follows that $\bar{g}_1 g_2 \in L_q^q L_\infty^p \subset L_q^1$. Moreover the finiteness of (10) shows that $\bar{g}_1 g_2 \in H_q^1$, and it clearly annihilates analytic functions. Applying the same functional analysis arguments as before, we conclude that

$$(11) \quad \begin{aligned} & (VMO_p/\mathcal{B}_0)^* \\ &= \mathcal{B}_0^\perp = \left\{ \sum_j \bar{h}_j g_j : h_j \in A^p, g_j \in (A^p)^\perp, \sum \|h_j\|_p \|g_j\|_q < \infty \right\}. \end{aligned}$$

Moreover, this is also the predual of BMO_p/\mathcal{B} . This gives us half of the first equivalence in the following theorem. Let A_0^p denote the class of functions f in A^p with $\int_D f dV = 0$.

4.8. Theorem. *The following are equivalent for a function f in L_q^1 :*

- (a) $f \in H_q^1$, that is, f is an l^1 sum of q -atoms.
- (b) $f = h + \sum_j \bar{h}_j g_j$ with $h \in A_0^1$, each $h_j \in A^p$, each $g_j \in (A^p)^\perp$ and $\sum \|h_j\|_p \|g_j\|_q < \infty$.
- (c) In case $n = 1$ and $D = \mathbb{D}$, $Pf \in L^1$ and $T^*f \in L^1$.

The decomposition in (b) can be associated with a norm, namely the infimum of $\|h\|_1 + \sum_j \|h_j\|_p \|g_j\|_q$ over all such decompositions. This norm is equivalent to the atomic norm. Similarly the sum $\|f\|_{L_q^1} + \|Pf\|_1 + \|T^*f\|_1$ is an equivalent norm. The operator T^* is the formal adjoint of T , that is, the operator with kernel $T^*(z, w) = \overline{T(w, z)}$.

Proof. We have seen that P is bounded from VMO_p to \mathcal{B}_0 , and we will see (in §5) that P is also bounded from H_q^1 to L^1 . Thus if $f \in H_q^1$ then we have $f = Pf + (I - P)f$. The second term belongs to \mathcal{B}_0^\perp and the first to A_0^1 . Then (b) follows from (a) by the preceding discussion. Conversely, let f have the form in (b). We need only show that $A_0^1 \subset H_q^1$ to conclude that $f \in H_q^1$. But given $h \in A_0^1$ and $g \in C(\bar{D})$ we have

$$\int_D h \bar{g} dV = \int_D Ph \bar{g} dV = \int_D h \overline{Pg} dv.$$

It is known that A_0^1 is the dual of \mathcal{B}_0 , and so

$$\left| \int_D h \overline{Pg} dV \right| \leq C \|h\|_1 \|Pg\|_{\mathcal{B}} \leq C \|h\|_1 \|g\|_{BMO, p}.$$

We conclude that as a linear functional on VMO_p , h agrees with an element of H_q^1 . It is easy to see that this means they are the same function.

Now let $n = 1$ and assume (a). We have already mentioned that $Pf \in L^1$ when $f \in H_q^1$, and it is similarly easy to show that T^* is bounded from atoms to L^1 . This gives the implication (a) \Rightarrow (c).

Now suppose that f is a function in L_q^1 such that Pf and T^*f are in L^1 . To show that f is in H_q^1 it suffices to show that it defines a bounded linear functional on VMO_p . To this end let $g \in VMO_p$ with $g = g_1 + T\varphi + P\psi$, where g_1, φ and ψ have compact support, and observe that

$$\begin{aligned} \langle f, g \rangle &= \langle f, g_1 \rangle + \langle f, T\varphi \rangle + \langle f, P\psi \rangle \\ &= \langle f, g_1 \rangle + \langle T^*f, \varphi \rangle + \langle Pf, \psi \rangle. \end{aligned}$$

It follows easily that

$$\begin{aligned} |\langle f, g \rangle| &\leq \|f\|_{L_q^1} \|g_1\|_{L_p^\infty} + \|T^*f\|_1 \|\varphi\|_\infty + \|Pf\|_1 \|\psi\|_\infty \\ &\leq (\|f\|_{L_q^1} + \|T^*f\|_1 + \|Pf\|_1) \|g\|_{BMO, p}. \end{aligned}$$

Since functions with compact support are dense in V_p^∞ as well as C_0 , this defines f as a bounded linear functional on a dense subset of VMO_p , and so (c) \Rightarrow (a). \square

Remark. It is a little more complicated to state a strongly pseudoconvex domain version of part (c) of Theorem 4.8, but here is an indication of how to go about it for the interested reader: Use the sup norm of the expression

$$\sup_{z \in D} (|\rho(z)\varphi(z)| + |\rho(z)|^{1/2} |\bar{\partial}\rho(z) \wedge \varphi(z)|)$$

to define an L^∞ -type space of forms. Its predual will be an appropriate L^1 -type space of forms. In place of T use the operator $\varphi \mapsto \int \varphi \wedge L(\cdot, w)$ from §6 (call it simply L), and then in place of (c) will be a similar statement involving L^* and that L^1 -type space.

5. BOUNDEDNESS OF HANKEL OPERATORS IN L^1 NORM

Since A_0^1 is contained in H_q^1 and the operators T , T^* , and P are bounded on H_q^1 it makes sense to consider the boundedness of various operators on A^1 defined using these operators. Here we investigate the boundedness of Hankel operators from A^1 to L^1 . The crucial problem, which makes this case different from the rest, is this: multiplying by a bounded function is not in general bounded from A^1 to H_q^1 for any q . It is bounded from A^1 to L^1 , but then neither P nor T is bounded on L^1 . We do not at present have a complete characterization of the Hankel operators which are bounded in L^1 norm. We present here one necessary condition and one sufficient condition. We then show that for the important special case of the conjugate of an analytic function, the sufficient condition is also necessary. Finally, we present an example to show that, in general, the sufficient condition is not necessary (and a simple argument to show that the necessary condition is not sufficient).

The necessary condition we obtain is gotten by following the proof of the corresponding result for A^p with $p > 1$, but refraining from continuing to the last step (d) \Rightarrow (a), which will not work.

5.1. Theorem. *Let f be a measurable function in L^1 and let H_f be defined on bounded analytic functions. If H_f extends to a bounded operator from A^1 to L^1 then the following equivalent conditions hold.*

(a) *For some (respectively, every) $r > 0$,*

$$F(z) \stackrel{\text{def}}{=} \inf_{h \in \mathcal{H}} \frac{1}{|B(z, r)|} \int_{B(z, r)} |f - h| dV$$

is bounded on d .

(b) *$f = f_1 + f_2$, where*

$$G(z) \stackrel{\text{def}}{=} \frac{1}{|B(z, r)|} \int_{B(z, r)} |f_1| dV$$

is bounded on D for some (respectively, every) $r > 0$ and

$$|\rho(z)| |\bar{\partial} f_2(z)| + |\rho(z)|^{1/2} |\bar{\partial} f_2(z) \wedge \partial \rho(z)|$$

is bounded on D .

(c) *$f \in BMO_1 + \mathcal{H}$.*

In particular, if f is the conjugate of an analytic function $f = \bar{h}$ then it follows from (a), (b), or (c) that h is in \mathcal{B} . The proof of Theorem 5.1 is

omitted, as it is identical to the proofs in [Lu3] and [Li2]. (That (c) implies (a) is clear, so the conditions are indeed equivalent.)

To obtain a sufficient condition we can try to find additional conditions on f_1 and f_2 which will make both H_{f_1} and H_{f_2} bounded, and then impose conditions on f that will produce these. Since the necessary condition on f_1 already implies M_{f_1} is bounded from A^1 to L^1 , the boundedness of PM_{f_1} will imply the boundedness of H_{f_1} . A sufficient condition for the boundedness of PM_{f_1} has already been obtained by K. Zhu for the ball, and this leads to the sufficient condition we consider in the next theorem.

5.2. Theorem. *Let f be a measurable, locally integrable function on D . Suppose that the following function is bounded:*

$$(12) \quad z \mapsto \log \frac{1}{|\rho(z)|} \inf_{h \in \mathcal{H}} \frac{1}{|B(z)|} \int_{B(z)} |f - h| dV.$$

Then there exists a decomposition $f = f_1 + f_2$, where

$$\log \frac{1}{|\rho(z)|} \frac{1}{|B(z, r)|} \int_{B(z, r)} |f_1| dV$$

is bounded on D for some (respectively, all) $r > 0$ and

$$(13) \quad \log \frac{1}{|\rho(z)|} [|\rho(z)| |\bar{\partial} f_2(z)| + |\rho(z)|^{1/2} |\bar{\partial} f_2(z) \wedge \partial \rho(z)|]$$

is bounded on D . Moreover, this decomposition implies that H_f extends to a bounded operator from A^1 to L^1 .

Proof. The methods of [Lu3] and [Li2] show that the boundedness of (12) implies the decomposition in the same way that (b) implies (c) in Theorem B and in the previous theorem.

We need to establish a series of lemmas. The difficulty (at least for the general strongly pseudoconvex domain) is that we do not have an explicit integral formula for $H_{f_2}g(z)$. This makes us follow a somewhat convoluted path.

We begin with a construction from [Li2] of a kernel $L(z, w)$ which will solve the $\bar{\partial}$ -equation with good estimates. Recall that ρ is the strictly plurisubharmonic defining function for D . We use $\rho_j(z)$ for the derivative $\partial \rho(z)/\partial z_j$ and ρ_{jk} for the corresponding second derivative, $1 \leq j, k \leq n$. Let $F(z, w)$ denote the Levi polynomial for D . That is,

$$F(z, w) = \sum_{j=1}^n \rho_j(w)(w_j - z_j) - \frac{1}{2} \sum_{j=1}^n \sum_{k=1}^n \rho_{jk}(w)(w_j - z_j)(w_k - z_k).$$

It is well known that there exist constants δ and c such that F satisfies

$$\operatorname{Re} F(z, w) + |\rho(w)| \geq \frac{|\rho(z)| + |\rho(w)|}{2} + c|z - w|^2$$

for all $z, w \in \bar{D}$ with $|z - w| < \delta$. The following lemma is from [F].

5.3. Lemma. *Let ρ and δ be as above. There exist functions $h_j(z, w)$, $1 \leq j \leq n$, and $\Psi(z, w)$ in $C^\infty(\bar{D} \times \bar{D})$ and a constant $C > 0$ with the following properties.*

(a) *For each fixed $w \in \bar{D}$, $h_j(z, w)$ and $\Psi(\cdot, w)$ are holomorphic on \bar{D} .*

- (b) $\Psi(z, w) = |\rho(w)| + \sum_{j=1}^n h_j(z, w)(w_j - z_j)$.
 (c) $|\Psi(z, w)| \geq 1/C$ when $|z - w| > \delta/2$, and otherwise $\Psi(z, w) = \frac{g(z, w)(F(z, w) + |\rho(w)|)}{\bar{D} \times \bar{D}}$, where g is smooth and nonvanishing on $\bar{D} \times \bar{D}$.
 (d) $h_j(w, w) = \rho_j(w) - \rho(w)g_j(w)$, where $g_j(w) = \partial g(z, w)/\partial w_j|_{z=w}$.

Now let

$$s(z, w) = \sum_{j=1}^n (\bar{w}_j - \bar{z}_j) dw_j, \quad h(z, w) = \sum_{j=1}^n h_j(z, w) dw_j,$$

$$q(z, w) = \frac{\bar{\partial}_w h(z, w)}{\rho(w)} - \frac{\bar{\partial} \rho(w) \wedge h(z, w)}{\rho(w)^2},$$

and finally for any $N \geq 1$ let

$$L^{(N)}(z, w) = C_{n, N} \sum_{k=0}^{n-1} a_k \left(\frac{\rho(w)}{\Psi(z, w)} \right)^{k+N} \frac{s(z, w) \wedge q(z, w)^k \wedge (\bar{\partial}_w s(z, w))^{n-k-1}}{|z - w|^{2n-2k}}$$

where a_k and $C_{n, N}$ are constants whose exact value need not concern us. This formula and the lemma to follow may be found in [Li2] and ultimately in [BA, p. 102, equation (20)].

5.4. Lemma. Let $N \geq 1$, $f \in C^1 \cap L^1$ on D and $|\rho \bar{\partial} f| + |\rho|^{1/2} |\bar{\partial} \rho \wedge \bar{\partial} f| \in L^1$. Then, for appropriate choices of $C_{n, N}$ and a_k ,

$$f(z) = \int_D \bar{\partial} f(w) \wedge L^{(N)}(z, w) + a_n \int_D f(w) \left(\frac{\rho(w)}{\Psi(z, w)} \right)^{n+N} \wedge q(z, w)^n,$$

where a_n is another constant whose value will not concern us. Moreover, if φ is a $\bar{\partial}$ -closed $(0, 1)$ -form with $|\rho \varphi| + |\rho|^{1/2} |\varphi \wedge \bar{\partial} \rho| \in L^1$, then

$$u(z) = \int_D \varphi(w) \wedge L^{(N)}(z, w)$$

solves the $\bar{\partial}$ -equation $\bar{\partial} u = \varphi$ in D .

Let $L_k(z, w)$ denote the k th term of the function for $L^{(N)}$ (omitting the constants). Thus

$$L_k(z, w) = \left(\frac{\rho(w)}{\Psi(z, w)} \right)^{k+N} \frac{s(z, w) \wedge q(z, w)^k \wedge (\bar{\partial}_w s(z, w))^{n-k-1}}{|z - w|^{2n-2k}}.$$

For the purposes of our proof, let us remark that

$$(14) \quad q(z, w)^k = \frac{(\rho \bar{\partial}_w h)^k - k(\rho \bar{\partial}_w h)^{k-1} \wedge h \wedge \bar{\partial} \rho}{\rho^{2k}} \\ = \rho^{-k-1} (\bar{\partial}_w h)^{k-1} \wedge (\rho \bar{\partial}_w h - k h \wedge \bar{\partial} \rho).$$

This is easily established by induction using $\bar{\partial} \rho \wedge \bar{\partial} \rho = 0$. Thus

$$(15) \quad L_k(z, w) = \frac{\rho(w)^{N-1} s(z, w) \wedge [\bar{\partial}_w s(z, w)]^{n-k-1}}{\Psi(z, w)^{k+N} |z - w|^{2n-2k}} \\ \wedge [\rho(w) q_1(z, w) + q_2(z, w) \wedge \bar{\partial} \rho(w)]$$

where q_1 and q_2 are both smooth and bounded. When $k = 0$ we have $q_1 = 1$ and $q_2 = 0$. Thus we have the following estimates if φ satisfies the condition of Lemma 3.1, namely $|\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho|$ is bounded:

$$\begin{aligned} |\varphi(w) \wedge L_0(z, w)| &\leq C \frac{\rho(w)^{N-1}}{|\Psi(z, w)|^N |z - w|^{2n-1}}, \\ |\varphi(w) \wedge L_1(z, w)| &\leq C \frac{\rho(w)^{N-3/2}}{|\Psi(z, w)|^{1+N} |z - w|^{2n-3}}, \\ |\varphi(w) \wedge L_k(z, w)| &\leq C \frac{\rho(w)^{N-3/2}}{|\Psi(z, w)|^{3/2+N} |z - w|^{2n-4}} \quad \text{if } k \geq 2, \end{aligned}$$

where the constant C is a multiple of the L^∞ norm of $|\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho|$. The last inequality uses the fact that $|\Psi(z, w)| \geq c|z - w|^2$.

Now let $K^{(N)}(z, w)$ be the kernel in the second integral of the formula for $f(z)$ in Lemma 5.4, that is,

$$K^{(N)}(z, w) dV(w) = \left(\frac{|\rho(w)|}{|\Psi(z, w)|} \right)^{n+N} q(z, w)^n,$$

where q and Ψ are as in Lemma 5.3. Let $H(z, w) = K^{(1)}(z, w)$. Let $A(z, w) = \overline{H(w, z)} - H(z, w)$. Let H and A denote the integral operators with kernels $H(z, w)$ and $A(z, w)$, respectively.

5.5. Lemma. *The Bergman projection P may be written $P = H + PA$. Moreover, $|A(z, w)| \leq C|\Psi(z, w)|^{-n-1/2}$, and $|\nabla_w A(z, w)| \leq C|\Psi(z, w)|^{-n-3/2}$.*

Proof. See [R, Chapter VII, §7]. \square

5.6. Lemma. *Let $t \geq 0$. The function Ψ satisfies the following estimates for all $z \in D$:*

$$\int_D \frac{|\rho(w)|^t}{|\Psi(z, w)|^{n+1+s+t}} + \frac{|\rho(w)|^t}{|\Psi(w, z)|^{n+1+s+t}} dV(w) \leq C \begin{cases} 1 & \text{if } s < 0, \\ |\log |\rho(z)|| & \text{if } s = 0, \\ |\rho(z)|^{-s} & \text{if } s > 0, \end{cases}$$

and if φ is a form satisfying the condition of Lemma 3.1 and if $N \geq 2$, then

$$(16) \quad \int_D |\varphi(w) \wedge L^{(N)}(z, w)| dV(z) \leq C \| |\rho\varphi| + |\rho|^{1/2}|\varphi \wedge \bar{\partial}\rho| \|_\infty.$$

Proof. The first part is almost [R, Lemma 7.3, p. 310]. The proof there relies only on the fact that $|\Psi(z, w)| \geq c(|\rho(z)| + |\rho(w)| + |z - w|^2)$ on $\bar{D} \times \bar{D}$. Thus it is valid with the variables interchanged. That lemma is stated with $t = 0$, but since $|\rho(w)|/|\Psi(z, w)| \leq C$ for all z and w , it clearly implies our statement.

The last estimate can be obtained by first using the appropriate estimate from the list preceding Lemma 5.5, then using the same lower estimate on $|\Psi|$ as above, and making a standard change of variables $z \rightarrow t(z)$ with $\text{Re } t_1(z) = \rho(w) - \rho(z)$ as in [Li2] or [DH]. This leads to one of the integrals in [Li2, Lemma 3.2] or [DH, Lemma 2.2]. The estimates there (with $\alpha = 0$) yield the estimate (16). \square

The main estimates we need are contained in the following lemma.

5.7. Lemma. *If $f \in C^1$ on D with $\log(1/|\rho|)(|\rho\bar{\partial}f| + |\rho|^{1/2}|\bar{\partial}f \wedge \bar{\partial}\rho|) \in L^\infty$ and $|f| \leq C(1 + |\log(|\rho|)|)$, then the Hankel operator H_f is bounded from A^1 to L^1 .*

Proof. Note that $H_f = (I - P)M_f = (I - H)M_f - PAM_f$. We estimate these two parts separately. Let g be an analytic function, continuous on \bar{D} . Then by Lemma 5.4 with $N = 1$ applied to g

$$\begin{aligned} (I - H)(fg)(z) &= \int (f(z) - f(w))g(w)H(z, w) dV(w) \\ &= C \int (f(z) - f(w))g(w) \\ &\quad \times \frac{[\rho(w)\bar{\partial}_w h(z, w) - n\bar{\partial}\rho(w) \wedge h(z, w)] \wedge \bar{\partial}_w h(z, w)^{n-1}}{\Psi(z, w)^{n+1}} \end{aligned}$$

from equation (14). Continuing the above equations:

$$\begin{aligned} (I - H)(fg)(z) &= Cf(z) \int g(w) \frac{\rho(w)(\bar{\partial}_w h(z, w))^n}{\Psi(z, w)^{n+1}} \\ &\quad - C \int f(w)g(w) \frac{\rho(w)(\bar{\partial}_w h(z, w))^n}{\Psi(z, w)^{n+1}} \\ &\quad - nC \int (f(z) - f(w))g(w) \frac{\bar{\partial}\rho(w) \wedge h(z, w) \wedge (\bar{\partial}_w h(z, w))^{n-1}}{\Psi(z, w)^{n+1}} \\ &= I_1(z) + I_2(z) + I_3(z). \end{aligned}$$

For any $\alpha > 1$ with $\alpha(n + 1/2) < n + 1$ one has from Lemma 5.6

$$(17) \quad \int_D \left(\frac{|\rho(w)|^{1/2}}{|\Psi(z, w)|^{n+1}} \right)^\alpha dV(w) + \int_D \left(\frac{|\rho(w)|^{1/2}}{|\Psi(z, w)|^{n+1}} \right)^\alpha dV(z) \leq C.$$

Note also that $|f(z)| \leq C(1 + |\log(|\rho|)|)$ implies that $f \in L^p$ for all $p > 0$ and $|\rho(z)|^{1/2}f(z) \in L^\infty$. The extension of Young's theorem [R, p. 358] together with (17) implies that the kernel

$$\frac{|\rho(w)|^{1/2}\bar{\partial}h(z, w)^n}{\Psi(z, w)^{n+1}}$$

is a bounded operator from L^1 to L^α . This implies that $\|I_1\|_1 + \|I_2\|_1 \leq C\|g\|_1$.

To estimate I_3 we use integration by parts. Recall that $\Psi(z, w) = -\rho(w) + \sum_{i=1}^n h_i(z, w)(w_i - z_i)$. By Stokes' theorem,

$$\begin{aligned} I_3(z) &= C \left[\int_D -\bar{\partial}f(w)g(w) \frac{\rho(w) \wedge h(z, w) \wedge \bar{\partial}_w h(z, w)^{n-1}}{\Psi(z, w)^{n+1}} \right. \\ &\quad + \int_D (f(z) - f(w))g(w) \frac{\rho(w) \wedge (\Psi(z, w)\bar{\partial}_w h(z, w)^n + H_1(z, w))}{\Psi(z, w)^{n+2}} \\ &\quad \left. + (n+1) \int_D (f(z) - f(w))g(w) \frac{\rho(w)\bar{\partial}\rho(w) \wedge h(z, w) \wedge \bar{\partial}_w h(z, w)^{n-1}}{\Psi(z, w)^{n+2}} \right] \\ &= J_1(z) + J_2(z) + J_3(z), \end{aligned}$$

where H_1 is an (n, n) -form with $|H_1(z, w)| \leq C|z - w|$. By Lemma 5.6, one can easily estimate that

$$\|J_1\|_1 \leq C\|(\log|\rho|)\rho\bar{\partial}f\|_\infty\|g\|_1.$$

Now observe that

$$|\rho(w)|^{1/2} \frac{|\Psi(z, w)(\bar{\partial}h)^n + H_1(z, w)|}{|\Psi(z, w)|^{n+2}} \leq \frac{C}{|\Psi(z, w)|^{n+1}}.$$

Then the same reasoning as in the estimates of $\|I_1\|_1$ and $\|I_2\|_1$ yields

$$\|J_2\|_1 \leq C\|g\|_1.$$

By Lemma 5.4, it is easy to see that there are constants C_1 and C_2 such that

$$\begin{aligned} J_3(z) &= C_1 \left[f(z)g(z) - \int_D f(w)g(w)K^{(2)}(z, w) dV(w) \right] \\ &\quad + C_2 \int_D (f(z) - f(w))g(w) \frac{\rho(w)^2 \wedge (\bar{\partial}h(z, w))^n}{\Psi(z, w)^{n+2}} \\ (18) \quad &= C_1 \int_D L^{(2)}(z, w) \wedge \bar{\partial}f(w)g(w) \\ &\quad + C_2 \int_D (f(z) - f(w))g(w) \frac{\rho(w)^2 \wedge (\bar{\partial}h(z, w))^n}{\Psi(z, w)^{n+2}}. \end{aligned}$$

By the same arguments as those used in the estimates of $\|I_1\|_1$ and $\|I_2\|_1$, we have that the L^1 -norm of the second term in (18) is bounded by $C\|g\|_1$.

By the second part of Lemma 5.6

$$\int_D |L^{(2)}(z, w) \wedge \bar{\partial}f(w)| dV(z) \leq C\|\rho\bar{\partial}f\| + |\rho|^{1/2}\|\bar{\partial}f \wedge \bar{\partial}\rho\|_\infty.$$

Thus

$$\|J_3\|_1 \leq C \int_D \int_D |L^{(2)}(z, w) \wedge \bar{\partial}f(w)g(w)| dV(z) dV(w) \leq C\|g\|_1.$$

Thus we have shown that $\|(I - H)M_f g\|_1 \leq C\|g\|_1$ for $g \in \mathcal{H} \cap C(\bar{D})$. Since such functions are dense in A^1 it follows that $(I - H)M_f$ is bounded from A^1 to L^1 . We turn now to the boundedness of PAM_f .

Let $\{U_i : 1 \leq i \leq m\}$ be an open cover of \bar{D} such that $\bar{U}_m \subset D$, $\partial D \subset \bigcup_{i=1}^{m-1} U_i$, and such that in each U_k we have $|\partial\rho/\partial z_{n_k}| > b > 0$ for some $n_k \leq n$, where b is a constant depending only on ρ . Let $\{\gamma_j\}$ be a smooth partition of unity subordinate to $\{U_i\}$. Then, by Stokes' theorem,

$$\begin{aligned} A(fg)(z) &= \sum_{k=1}^{m-1} \int_D A(z, w) \frac{\gamma_k(w)}{\partial\rho/\partial\bar{w}_{n_k}} f(w)g(w) \frac{\partial\rho}{\partial\bar{w}_{n_k}} dV(w) \\ &\quad + \int_D \gamma_m(w) f(w)g(w) A(z, w) dV(w) \\ &= \sum_{k=1}^{m-1} \int_D \left[F_k(w) \frac{\partial A(z, w)}{\partial\bar{w}_{n_k}} + G_k(w) A(z, w) \right] \rho(w) f(w)g(w) dV(w) \\ &\quad + \sum_{k=1}^{m-1} \int_D F_k(z, w) A(z, w) \rho(w) \frac{\partial f}{\partial\bar{w}_{n_k}} g(w) dV(w) \\ &\quad + \int_D \gamma_m(w) f(w)g(w) A(z, w) dV(w), \end{aligned}$$

where $F_k = \gamma_k / (\partial \rho / \partial \bar{w}_{n_k})$ and $G_k = \partial F_k / \partial \bar{w}_{n_k}$ are all smooth on \bar{D} . The assumptions imply that $\gamma_m f$ is bounded, so we get

$$|A(fg)(z)| \leq C \int_D \frac{|\rho(w)f(w)g(w)|}{|\Psi(z, w)|^{n+3/2}} + \frac{|\rho(w)\bar{\partial}f(w)g(w)| + |g(w)|}{|\Psi(z, w)|^{n+1/2}} dV(w).$$

The hypotheses on f imply that $|\rho|^{1/4}f$ and $\rho\bar{\partial}f$ are bounded. Since Ψ and ρ/Ψ are also bounded, we get

$$|A(fg)(z)| \leq C \int_D \frac{|g(w)|}{|\Psi(z, w)|^{n+3/4}} dV(w).$$

For any $\alpha > 1$ with $\alpha(n + 3/4) < n + 1$ we get from Lemma 5.6

$$\int_D |\Psi(z, w)^{-n-3/4}|^\alpha dV(w) + \int_D |\Psi(z, w)^{-n-3/4}|^\alpha dV(z) \leq C.$$

Then [R, p. 358] and the preceding inequalities imply that $\|A(fg)\|_\alpha \leq C\|g\|_1$. Since P is bounded on L^α for any $\alpha > 1$, we have $\|PA(fg)\|_1 \leq C\|PA(fg)\|_\alpha \leq C\|g\|_1$, and the proof of Lemma 5.7 is complete. \square

Proof (of Theorem 5.2). The only thing that prevents us from using Lemma 5.7 directly is the hypothesis that $|f(z)| \leq C(1 + |\log|\rho(z)||)$ and the fact that the constants depended on this C . To get around this we write $f = f_1 + f_2$ as in the hypothesis. Let $u(z)$ be the solution to $\bar{\partial}u = \bar{\partial}f_2$ guaranteed by Lemma 3.1, and observe two things: (1) Since $\bar{\partial}u = \bar{\partial}f_2$, u satisfies the same condition as f_2 . (2) Since f_2 satisfies (13), it satisfies the weaker condition

$$|\rho(z)\bar{\partial}f_2(z)| + |\rho(z)|^{1/2}|\bar{\partial}f_2(z) \wedge \partial\rho(z)| \in L^\infty,$$

and this is sufficient by Lemma 3.1 to give $|u(z) - u(w)| \leq C(1 + \beta(z, w))$ with the constant C depending only on the L^∞ norm of the above expression. This easily implies

$$|u(z)| \leq C(1 + |\log|\rho(z)||).$$

Thus H_u is bounded by Lemma 5.7, and so H_{f_2} is bounded since it equals H_u .

To finish, we examine H_{f_1} . Again the condition on f_1 implies a weaker condition:

$$(19) \quad \frac{1}{|B(z, r)|} \int_{B(z, r)} |f_1| dV \in L^\infty.$$

This implies that M_{f_1} is bounded from A^1 to L^1 , and it remains to show that PM_{f_1} is also bounded. Now $PM_{f_1}g(z) = \int_D K(z, w)f_1(w)g(w)dV(w)$ and $\int_D |K(z, w)|dV(w) \leq C(1 + |\log|\rho(w)||)$, which implies that

$$(20) \quad \|PM_{f_1}g\|_1 \leq C \int_D (1 + |\log|\rho||)|f_1||g|dV.$$

Letting $f_3 = (1 + |\log|\rho||)f_1$ and observing that $\rho(z) \sim \rho(w)$ when $z, w \in B(z, r)$, we conclude that f_3 also satisfies the condition (19). Thus f_3 is a multiplier from A^1 to L^1 , and this means the right-hand side of (20) is bounded by a multiple of the L^1 norm of g . \square

Our next result is to show the converse of Theorem 5.2 when the symbol f is the conjugate of an analytic function.

Recall the definition of q -atom from §4. These are functions a which satisfy the following conditions (remembering $r > 0$ is fixed):

a is supported in some ball $B(z, r)$,

$$\int_D a(w) dV(w) = 0,$$

$$\|a\|_q = |B(z, r)|^{1/q-1}.$$

The space H_q^1 consists of all functions of the form $f(z) = \sum \lambda_j a_j(z)$ where the a_j are q -atoms and $\sum |\lambda_j| < \infty$. The norm of f in H_q^1 is the infimum of $\sum |\lambda_j|$ over all such representations of f . Note that all atoms satisfy $\|a\|_1 \leq 1$.

5.8. Lemma. For $q \geq 1$, the Bergman projection P is bounded from H_q^1 to the Bergman space A^1 .

Proof. It suffices to show that $\|Pa\|_1 \leq C$ for all atoms a with a constant C independent of a . Let a be an atom supported in $B(\lambda, r)$. Because of the cancellation property of an atom, we have

$$Pa(z) = \int_D K(z, w) a(w) dV(w) = \int_D [K(z, w) - K(z, \lambda)] a(w) dV(w).$$

Note that for $w \in B(\lambda, r)$,

$$|\rho(w)| \sim |\rho(z)|, \quad |w - \lambda| \leq C|\rho(\lambda)|^{1/2},$$

$$\left| \sum \frac{\partial \rho(\lambda)}{\partial z_j} (w_j - \lambda_j) \right| \leq C|\rho(\lambda)|.$$

A direct calculation yields that

$$|\Psi(z, w) - \Psi(z, \lambda)| \leq C[|\rho(\lambda)|^{1/2}(|z - w| + |z - \lambda|) + |\rho(\lambda)|]$$

$$\leq C[|\rho(\lambda)|^{1/2}(|\Psi(z, w)|^{1/2} + |\Psi(z, \lambda)|^{1/2}) + |\rho(\lambda)|].$$

By C. Fefferman's formula for the Bergman kernel [Fe], it follows that

$$|K(z, w) - K(z, \lambda)| \leq C \left[\frac{|\rho(\lambda)|}{|\Psi(z, \lambda)|^{n+2}} + \frac{|\rho(\lambda)|^{1/2}}{|\Psi(z, \lambda)|^{n+3/2}} \right]$$

$$+ C \left[\frac{|\rho(w)|}{|\Psi(z, w)|^{n+2}} + \frac{|\rho(w)|^{1/2}}{|\Psi(z, w)|^{n+3/2}} \right]$$

$$\leq C \left[\frac{|\rho(\lambda)|^{1/2}}{|\Psi(z, \lambda)|^{n+3/2}} + \frac{|\rho(w)|^{1/2}}{|\Psi(z, w)|^{n+3/2}} \right].$$

By Fubini's theorem and Lemma 5.6 we have

$$\|Pa\|_1 \leq C\|a\|_1 \int_D \frac{|\rho(\lambda)|^{1/2}}{|\Psi(z, \lambda)|^{n+3/2}} dV(z) \leq C',$$

where C and C' are constants independent of a . \square

5.9. Lemma. Let $f \in A^1$. If the Hankel operator $H_{\bar{f}}$ is bounded from A^1 to L^1 , then f is in the Bloch space \mathcal{B} and $|f(z) - f(w)| \leq C\beta(z, w)$.

Proof. We have already seen that the boundedness of $H_{\bar{f}}$ implies that $f \in BDA_1$ (Theorem 5.1). That is to say, there is a constant C such that for any

ball $B = B(z, r)$ there is an analytic function h such that $\int_B |\bar{f} - h| dV \leq C|B|$. Now the same will be true if B is replaced with a polydisk with radius $c|\rho(z)|$ in the complex direction normal to the boundary and radii $c\sqrt{|\rho(z)|}$ in the complex directions parallel to the boundary. It is straightforward to extend the inequality $\int_{\mathbb{D}} |\bar{f} - \bar{f}(0)| dA \leq C \int_{\mathbb{D}} |\bar{f} - h| dA$ from the disk to a polydisk. (That is, a constant is nearly the best analytic approximant to \bar{f} in L^1 .) Thus $f \in BMO_1$. And we easily deduce from this (see [Li1] or even the remarks at the end of [Lu2]) that $f \in \mathcal{B}$ and that $|f(z) - f(w)| \leq C\beta(z, w)$. This gives our conclusion. \square

5.10. Theorem. *Let $f \in A^1$. Then $H_{\bar{f}}$ is bounded from A^1 to L^1 if and only if*

$$\left(\log \frac{1}{|\rho(z)|} \right) |\partial f(z)| \cdot |\rho(z)| \in L^\infty.$$

Proof. We need only prove the necessity. By Lemma 5.8, if $H_{\bar{f}}$ is bounded, then $H_{\bar{f}}P$ is bounded from H_2^1 to L^1 . Thus $\|H_{\bar{f}}Pa\|_1 \leq C$ for all atoms a . Writing $B = B(\lambda, r)$ for the Bergman metric ball that supports a (as in the definition of atom), it is easy to see that $H_{\bar{f}}Pa$ may be written

$$\begin{aligned} H_{\bar{f}}Pa(z) &= \int_D (\bar{f}(z) - \bar{f}(w))K(z, w)a(w) dV(w) \\ &= \int_B (\bar{f}(z) - \bar{f}_B)K(z, w)a(w) dV(w) \\ &\quad - \int_B (\bar{f}(w) - \bar{f}_B)K(z, w)a(w) dV(w) \\ &= I_1(z) + I_2(z). \end{aligned}$$

Note that we have the following for $I_1(z)$:

$$I_1(z) = (\bar{f}(z) - \bar{f}_B) \int_B (K(z, w) - K(z, \lambda))a(w) dV(w).$$

We claim that $\|I_1\|_1 \leq C$. By estimates in [Li1], there is a constant $\delta > 0$ such that when $|z - w| \leq \delta$ then for any ε with $0 < \varepsilon < 1/(4n + 4)$ we have

$$(21) \quad \beta(z, w) \leq C_\varepsilon (|\Psi(z, w)|^2 / |\rho(z)\rho(w)|)^\varepsilon.$$

Moreover we always have

$$(22) \quad \beta(z, w) \leq C[|\log |\rho(z)|| + |\log |\rho(w)||],$$

and when $|z - w| > \delta$ we have

$$(23) \quad |\Psi(z, w)| \geq 1/C > 0.$$

From Lemma 5.9 we have $|f(z) - f_B| \leq \beta + r$. From this and the proof of Lemma 5.8 we have

$$\|I_1\|_1 \leq C\|a\|_1 \int_D (\beta(z, \lambda) + r) \left[\frac{|\rho(\lambda)|^{1/2}}{|\Psi(z, \lambda)|^{n+3/2}} + \frac{|\rho(w)|^{1/2}}{|\Psi(z, w)|^{n+3/2}} \right] dV(z).$$

Finally, (21), (22), (23), $\beta(z, \lambda) \leq \beta(z, w) + r$, and the usual estimates (Lemma 5.6) show that $\|I_1\|_1 \leq C$.

Turning to I_2 , we write it as

$$\begin{aligned} I_2(z) &= \int_B (\bar{f}_B - \bar{f}(w))(K(z, w) - K(z, \lambda))a(w) dV(w) \\ &\quad + K(z, \lambda) \int_B (\bar{f}_B - \bar{f}(w))a(w) dV(w) \\ &= J_1(z) + J_2(z). \end{aligned}$$

Using the estimates above, we obtain

$$\|J_1\|_1 \leq C \|a\|_1 \int_D (\beta(w, \lambda) + r) \left[\frac{|\rho(\lambda)|^{1/2}}{|\Psi(z, \lambda)|^{n+3/2}} + \frac{|\rho(w)|^{1/2}}{|\Psi(z, w)|^{n+3/2}} \right] dV(z),$$

and again obtain $\|J_1\|_1 \leq C$. We are down to the following: Assuming $H_{\bar{f}}$ is bounded, we may conclude that $\|J_2\|_1 \leq C$ with a constant C depending only on f (and of course D). Thus

$$(24) \quad \|J_2\|_1 = \|K(\cdot, \lambda)\|_1 \left| \int_B (\bar{f}(w) - \bar{f}_B)a(w) dV(w) \right| \leq C$$

holds for all atoms a . We may suppose f is not constant. Now $f \in \mathcal{B}$ implies ([Li1]) that $0 < M_\lambda \stackrel{\text{def}}{=} |B|^{-1} \int_B |f(w) - f_B|^2 dV(w) \leq C$ for some constant C independent of λ . Thus $a_\lambda(z) \stackrel{\text{def}}{=} |B|^{-1} \chi_B(z)(f(z) - f_B)/M_\lambda^{1/2}$ is a 2-atom. Since $\|K(\cdot, \lambda)\|_1 \sim |\log |\rho(\lambda)||$, it follows from (24) that

$$|\log |\rho(\lambda)|| \left| \frac{1}{|B(\lambda, r)|} \int_{B(\lambda, r)} |f(w) - f_B|^2 dV(w) \right|^{1/2} \leq C.$$

It was proved in [Li1] that

$$|\rho(\lambda) \partial f(\lambda)| \leq C \left[\frac{1}{|B(\lambda, r)|} \int_{B(\lambda, r)} |f(w) - f_B|^2 dV(w) \right]^{1/2}.$$

Therefore $\log |\rho(z)| |\rho(z)| |\partial f(z)| \in L^\infty$. \square

If the necessary conditions of Theorem 5.1 were sufficient, this would contradict Theorem 5.10. The following example will show that the conditions of Theorem 5.2 are not necessary.

Example. Let us specialize to the unit disk \mathbb{D} , though the example is easily generalized to any domain D . Let S be a square centered at the origin and compactly contained in \mathbb{D} and let $M \geq 1$ be an integer. Divide S into $2M$ equal vertical stripes. Let ψ be a function which is constant on each stripe, takes there only the values ± 1 , does not take the same value on two stripes with a common boundary and is 0 off S . It is trivial to estimate for all $f \in C^1$

$$(25) \quad \left| \int_{\mathbb{D}} f \psi dA \right| \leq C \frac{\sup_{z \in S} |\nabla f(z)|}{M} \leq C_S \frac{\|f\|_1}{M}, \quad \text{if } f \in A^1$$

with the C and C_S independent of M . Moreover,

$$(26) \quad \int_S |\psi - h| dA > \delta$$

for all $h \in \mathcal{H}$ with δ independent of M and h .

Let $B_j = B(w_j, r)$ be a sequence of disjoint Bergman metric disks, and let M_j be a sequence of integers. Let S_j be the square inside B_j obtained as the image of S under the dilation and translation taking \mathbb{D} to B_j . Let ψ_j be defined on B_j in the same way that ψ is defined on \mathbb{D} , using the integer M_j for M . Finally, let φ denote the function which is zero off $\bigcup_j B_j$ and equal to ψ_j on each B_j . It is clear from (26) that φ does not satisfy the sufficient condition (b) of Theorem 5.2. In fact,

$$\frac{1}{|B_j|} \inf_{h \in \mathcal{H}} \int_{B_j} |\varphi - h| dA$$

does not tend to 0 and so $|\log(1 - |z_j|^2)|$ times it cannot be bounded. However, the following estimates show that PM_φ is bounded from A^1 to L^1 for some choice of the sequence M_j , and therefore $H_\varphi = M_\varphi - PM_\varphi$ is also bounded. For $g \in A^1$

$$\begin{aligned} |PM_\varphi g(z)| &\leq \sum \left| \int_{B_j} \psi_j(w) g(w) K(z, w) dA(w) \right| \\ &\leq C \sum r_j^3 \frac{\sup_{w \in S_j} |g'(w) K(z, w) + g(w) \partial_w K(z, w)|}{M_j} \\ &\leq C \sum \frac{1}{M_j} \int_{B_j} |g| dA [|K(z, w_j)| + r_j |\nabla_w K(z, w_j)|], \end{aligned}$$

by (25), where $r_j \sim (1 - |w_j|)$ is the Euclidean radius of B_j . Since $K(\cdot, w)$ and $\nabla_w K(\cdot, w)$ satisfy $\|K(\cdot, w) + (1 - |w|) |\nabla_w K(\cdot, w)|\|_1 \leq C(1 + |\log(1 - |w|)|)$, it suffices to choose $M_j \geq (1 + |\log(1 - |w_j|)|)$ to make $\|PM_\varphi g\|_1 \leq C\|g\|_1$.

6. THE PROOF OF LEMMA 3.1

Recall Lemma 5.3 from §5. We shall use it again to construct another kernel $L(z, w)$ for the solution of the $\bar{\partial}$ -equation. One difference between this new one and the $L^{(N)}(z, w)$ already constructed is that the factor $|z - w|^{2(n-k)}$ in the denominator is replaced by an expression $\|z - w\|_A^{2(n-k)}$ where $\|\cdot\|_A^2$ is a positive definite quadratic form whose behavior is better suited to the geometry of a strongly pseudoconvex domain (and to our estimates). Following [BA], this necessitates a different choice of the form s from the one defined prior to Lemma 5.4. Also, we will choose a slightly different change of variables from the one used in [Li2] and [DH]. The construction of L below first appeared in [Li3].

Let us write

$$G(z, w) = F(w, z) = \sum_{j=1}^n \rho_j(z)(z_j - w_j) - \frac{1}{2} \sum_{j=1}^n \sum_{k=1}^n \rho_{jk}(z)(z_j - w_j)(z_k - w_k)$$

where $F(z, w)$ is the Levi polynomial for D . Here are a few simple observations that we will need later.

$$\begin{aligned} |G(z, w) + F(z, w)| &= \left| \sum_{j=1}^n (\rho_j(w) - \rho_j(z))(w_j - z_j) + O(|z - w|^2) \right| \\ &\leq C|z - w|^2, \end{aligned}$$

and if $\delta > 0$ is as in Lemma 5.3 and $|z - w| < \delta$, then

$$\begin{aligned} |G(z, w)| &\leq C(|F(z, w)| + |z - w|^2) \\ &\leq C(|\operatorname{Im} F(z, w)| + |\operatorname{Re} F(z, w) + |\rho(w)|| + |\rho(w)| + |z - w|^2) \\ &\leq C(|\operatorname{Im} F(z, w)| + |\operatorname{Re} F(z, w) + |\rho(w)||) \leq C|\Psi(z, w)|. \end{aligned}$$

Similarly, $|F(z, w)| \leq C|\Psi(z, w)|$ and $|\rho(z)| + |\rho(w)| + |z - w|^2 \leq C|\Psi(z, w)|$ in $|z - w| < \delta$.

6.1. Lemma. *Let ρ and G be as above. Then there exist constants δ and C such that for any $z \in D$ with $|\rho(z)| \leq \delta$, in $|w - z| < \delta$ we can perform a smooth change of variables $t = t(w) = (t_1, t_2, \dots, t_n)$ with the following properties:*

- (a) $t_1(w) = G(z, w)$.
- (b) $|z - w|/C \leq |t(w)| \leq C|z - w|$ for w in $|w - z| < \delta$.
- (c) $1/C \leq |\partial t / \partial w| \leq C$ for w in $|w - z| < \delta$, where $\partial t / \partial w$ denotes the Jacobian of t .

Remark. It is standard (see [R] or [DH]) to change coordinates with a $\tau = \tau(w)$ where $\operatorname{Re} \tau_1(w) = \rho(w) - \rho(z)$ (or simply $\operatorname{Re} \tau_1(w) = \rho(w)$) and $\operatorname{Im} \tau_1(w) = \operatorname{Im} G(z, w)$. These have the advantage that level sets of ρ are mapped to hyperplanes, but the special nature of the integrals we will need to estimate seems to require us to use the slightly different coordinates of the lemma. (The real parts of t_1 and τ_1 differ only by $O(|z - w|^2)$ near z .)

Now choose a smooth function χ on $\mathbb{C}^n \times \mathbb{C}^n$ such that $0 \leq \chi \leq 1$ and

$$\chi(z, w) = \begin{cases} 1 & \text{if } |z - w| \leq \delta/2, \\ 0 & \text{if } |z - w| \geq \delta, \end{cases}$$

and write

$$\begin{aligned} G_1(z, w) &= \chi G(z, w) + (1 - \chi)|z - w|^2, \\ \|z - w\|_A^2 &= |G_1(z, w)|^2 + |\rho(z)||z - w|^2, \end{aligned}$$

and

$$\begin{aligned} s_j(z, w) &= \overline{G_1(z, w)} \left[\chi(z, w) \left(\rho_j(z) - \frac{1}{2} \sum_{k=1}^n \rho_{jk}(z_k - w_k) \right) \right. \\ &\quad \left. + (1 - \chi(z, w))(\overline{z_j - w_j}) \right] \\ &\quad + |\rho(z)|(\overline{z_j - w_j}). \end{aligned}$$

It is easy to calculate that $\|z - w\|_A^2 = \sum s_j(z, w)(z_j - w_j)$, that $|s_j| \leq C|z - w|$ uniformly for $z, w \in \overline{D}$, and that for z in any compact subset Ω of D , $\|z - w\|_A^2 \geq C'|z - w|^2$ uniformly for $w \in \overline{D}$. (The constant C' is allowed to depend on the compact set Ω .) Taken together, these statements say that the s_j satisfy condition (1) in [BA].

6.2. Lemma [Li3]. *The quantities G and $\|\cdot\|_A$ satisfy the following inequalities.*

- (a) *If $|z - w| < \delta/2$ then $\|z - w\|_A^2 \geq (1/C)(|G(z, w)|^2 + |\rho(w)||z - w|^2)$.*

(b) If $|z - w| \geq \delta/2$ then $\|z - w\|_A^2 \geq |\rho(w)|/C$.

Following B. Berndtsson and M. Andersson [BA] we define

$$s(z, w) = \sum_{j=1}^n s_j(z, w) dw_j, \quad h(z, w) = \sum_{j=1}^n h_j(z, w) dw_j,$$

$$q(z, w) = \bar{\partial}_w j(z, w)/\rho(w) - \bar{\partial}\rho(w) \wedge h(z, w)/\rho(w)^2,$$

and

$$L(z, w) = C_n \sum_0^{n-1} \gamma_k \left[\frac{|\rho(w)|}{\Psi(z, w)} \right]^{k+n+1} \frac{s(z, w) \wedge q(z, w)^k \wedge (\bar{\partial}_w s)^{n-k-1}}{\|z - w\|_A^{2(n-k)}},$$

where γ_k are constants defined in [BA, equation (20), p. 102] by letting $N = n + 1$ and $C_n = (-1)^{n(n+1)/2}/(n-1)!$. Also, h is the same h from Lemma 5.3 and q is the same as in §5.

6.3. Lemma [BA, pp. 103–104]. *If φ is a $\bar{\partial}$ -closed $(0, 1)$ form with coefficients in $C^1(\bar{D})$, then*

$$u(z) = \int_D \varphi \wedge L(z, w)$$

is a solution to the equation $\bar{\partial}u = \varphi$.

Remark. In [BA], this result was proved for strictly convex domains by letting

$$s_j(z, w) = \left[\sum_{k=1}^n \rho_k(z) (\overline{z_k - w_k}) \right] \rho_j(z) + |\rho(z)| (\overline{z_j - w_j})$$

and

$$h_j(z, w) = \rho_j(w),$$

for $1 \leq j \leq n$. As indicated in [BA, p. 104], an application of the same arguments yields the result stated here.

Remark. By a standard argument (e.g., see [DH, p. 453]), Lemma 6.3 may be extended to $\bar{\partial}$ -closed $(0, 1)$ forms with coefficients in $C^1(D) \cap L^1(D)$.

It is easy to check the following for $|z - w| < \delta/2$.

$$s(z, w) = \overline{G(z, w)} \left[\sum_1^n \rho_j(z) dw_j + \beta_1(z, w) \right] + |\rho(z)| \sum_1^n (\overline{z_j - w_j}) dw_j,$$

$$\bar{\partial}_w s(z, w) = \bar{\partial}_w \overline{G(z, w)} \wedge \left[\sum_1^n \rho_j(z) dw_j + \beta_1(z, w) \right] + \rho(z) \sum_1^n d\bar{w}_j \wedge dw_j$$

$$= \left[- \sum_1^n \bar{\rho}_j d\bar{w}_j + \overline{\beta_2(z, w)} \right] \wedge \left[\sum_1^n \rho_j(z) dw_j + \beta_1(z, w) \right]$$

$$+ \rho(z) \sum_1^n d\bar{w}_j \wedge dw_j,$$

where β_m , $m = 1, 2$, are $(1, 0)$ forms with $|\beta_m(z, w)| \leq C|z - w|$ and $\bar{\partial}_w \beta_1 = 0$.

Recall from §5 that

$$q(z, w)^k = \frac{\rho(w)(\bar{\partial}_w h(z, w))^k - k\bar{\partial}\rho(w) \wedge h(z, w) \wedge (\bar{\partial}_w h(z, w))^{k-1}}{\rho(w)^{k+1}}.$$

Also, for $|z - w| < \delta/2$ and $k < n - 1$,

$$(\bar{\partial}_w s)^{n-k-1} = [\rho a + (n - k - 1)\bar{\partial}_w \bar{G} \wedge \bar{s}] \wedge (\rho a)^{n-k-2}$$

where

$$a(z, w) = \sum_1^n d\bar{w}_j \wedge dw_j, \quad \bar{s}(z, w) = \sum_1^n \rho_j(z) dw_j + \beta_1(z, w).$$

It follows that we may write the combination $s(z, w) \wedge (\bar{\partial}_w s(z, w))^{n-k-1}$ that appears in the formula for $L(z, w)$ as

$$s \wedge (\bar{\partial}_w s)^{n-k-1} = \rho(z)^{n-k-1} \{ \bar{G} \bar{s} \wedge a + g_1 \wedge [\rho(z)a + (n - k - 1)\bar{\partial} \bar{G} \wedge \bar{s}] \} \wedge a^{n-k-2},$$

where $g_1(z, w) = \sum_1^n (\bar{z}_j - \bar{w}_j) dw_j$.

This gives us the following formula for $L(z, w)$:

$$L(z, w) = C_n \sum_{k=0}^{n-1} \gamma_k L_k(z, w)$$

with

$$\begin{aligned} L_k(z, w) &= \frac{\rho(z)^{n-k-1} \rho(w)^n}{\Psi(z, w)^{n+k+1} \|z - w\|_A^{2(n-k)}} [\rho(w)(\bar{\partial}_w h)^k - k\bar{\partial}\rho(w) \wedge h \wedge (\bar{\partial}_w h)^{k-1}] \\ &\quad \wedge \{ \bar{G} \bar{s} \wedge a + g_1 \wedge [\rho(z)a + (n - k - 1)\bar{\partial} \bar{G} \wedge \bar{s}] \} \wedge a^{n-k-2}, \end{aligned}$$

except that when $k = 0$ the first bracket becomes simply $\rho(w)$, and when $k = n - 1$ the entire expression on the last line becomes $\bar{G} \bar{s} + \rho g_1$ which equals $s(z, w)$.

For convenience we will use $\mathcal{E}_m(z, w)$ to denote smooth forms on $\bar{D} \times \bar{D}$ satisfying $|\mathcal{E}_m(z, w)| \leq C|z - w|^m$. The actual form represented by \mathcal{E}_m may change from one instance to another, even on the same line.

Note the following estimates:

$$\begin{aligned} \partial_w G_1(z, w) &= \partial_w G(z, w) = -\partial\rho(w) + \mathcal{E}_1, \\ \tilde{s}(z, w) &= \sum \rho_j(\cdot) dw_j + \beta_1(z, w) = \partial\rho(w) + \mathcal{E}_1, \\ h(z, w) &= \partial\rho(w) + \mathcal{E}_1 + \rho(w)\mathcal{E}_0. \end{aligned}$$

It follows that

$$\begin{aligned} \bar{\partial}_w \bar{G}_1(z, w) &= \bar{\partial}\rho(w) + \mathcal{E}_1(z, w), \\ \tilde{s}(z, w) \wedge h(z, w) &= \mathcal{E}_1(z, w) + \mathcal{E}_0(z, w)\rho(w), \\ \bar{\partial}_w \bar{G}_1 \wedge \tilde{s}(z, w) \wedge \bar{\partial}\rho(w) \wedge h(z, w) &= \bar{\partial}\rho(w) \wedge \mathcal{E}_1(z, w)[\mathcal{E}_1(z, w) + \rho(w)\mathcal{E}_0(z, w)]. \end{aligned}$$

Thus we have

$$(27) \quad L_k(z, w) = \frac{\rho(z)^{n-k-1} \rho(w)^n}{\Psi(z, w)^{n+k+1} \|z - w\|_A^{2(n-k)}} \{[(n-k-1)\mathcal{E}_2 + \rho(z)\mathcal{E}_1 + \bar{G}\mathcal{E}_0] \wedge \rho(w) + [k(\rho(z)\mathcal{E}_1 + \bar{G}\mathcal{E}_1) + (n-k-1)(\rho(w)\mathcal{E}_1 + k\mathcal{E}_3)] \wedge \bar{\partial}\rho(w)\}.$$

(We keep the k and the $n-k-1$ before certain terms rather than absorb them in the \mathcal{E}_m , because these terms vanish in the extreme cases $k=0$ or $k=n-1$.)

In the next lemma we give estimates for the solution of the $\bar{\partial}$ -equation. It seems to us that the standard integral representations and estimates [R, DH] do not work in our case. The following estimates may therefore be of some independent interest.

6.4. Lemma. *If φ is a $\bar{\partial}$ -closed $(0, 1)$ form with coefficients in $C^1(D) \cap L^1(D)$ and the expression $|\varphi \wedge \bar{\partial}\rho| |\rho|^{1/2} + |\rho\varphi|$ belongs to $L^\infty(D)$, then the solution*

$$(28) \quad u(z) = \int_D \varphi(w) \wedge L(z, w)$$

of the $\bar{\partial}$ -equation $\bar{\partial}u = \varphi$ has the form $u(z) = u_1(z) + u_2(z)$, where

$$\|u_1\|_\infty + \|\rho \nabla u_2\|_\infty \leq C \| |\rho\varphi| + |\rho|^{1/2} |\varphi \wedge \bar{\partial}\rho| \|_\infty$$

and

$$|\partial u_2(z) \wedge \partial \rho(z)| + |\bar{\partial} u_2(z) \wedge \bar{\partial} \rho(z)| \leq C |\rho(z)|^{-1/2} \| |\rho\varphi| + |\rho|^{1/2} |\varphi \wedge \bar{\partial}\rho| \|_\infty.$$

Proof. We may assume without loss of generality that the L^∞ norm of $|\varphi \wedge \bar{\partial}\rho| |\rho|^{1/2} + |\rho\varphi|$ is 1. By Lemmas 5.3 and (6.2), if $|z - w| \geq \delta/2$ then $\|\varphi \wedge L\|_\infty \leq C$. If $|\rho(z)| \geq \delta$ and $|z - w| < \delta$ then $|\varphi(w) \wedge L(z, w)| \leq C/|z - w|^{2n-1}$ and a similar estimate holds for $|\rho(w)| > \delta$. Thus the portion of the integral in (28) over $|w - z| > \delta/2$ is bounded independently of z and the portion of the integral over $|w - z| < \delta$ is bounded for $|\rho(z)| > \delta$. Thus, it suffices to obtain estimates when z is close to the boundary of D and we need only estimate the integral over $|w - z| < \delta/2$.

Now write

$$v_k(z) = \int_D \varphi(w) \wedge L_k(z, w).$$

We begin by showing that for $0 \leq k \leq n-2$, v_k is actually bounded. It is easy to check from (27) that for such k , $|\varphi(w) \wedge L_k(z, w)|$ is dominated by the sum of the two terms

$$\frac{|\rho(z)|^{n-k-1} |\rho(w)|^{n-1/2}}{\|z - w\|_A^{2(n-k)} |\Psi(z, w)|^{n+k+1}} (|G(z, w)| + |\rho(w)| + |\rho(z)| + |z - w|^2) |z - w|$$

and

$$\frac{|\rho(z)|^{n-k-1} |\rho(w)|^n}{\|z - w\|_A^{2(n-k)} |\Psi(z, w)|^{n+k+1}} (|G(z, w)| + |\rho(z)| |z - w| + |z - w|^2).$$

Note that, from Lemma 5.3 when $|z - w| \leq \delta$ and $|\rho(z)| < \delta$,

$$|\Psi(z, w)| \geq (\operatorname{Re} F(z, w) + |\rho(w)|)/C \geq (|\rho(z)| + |\rho(w)| + |z - w|^2)/C.$$

Recall also that $|\Psi(z, w)| \geq |G(z, w)|/C$. Thus the terms above are dominated by

$$\frac{|\rho(z)|^{n-k-1}|G(z, w)|}{\|z-w\|_A^{2(n-k)}|\Psi(z, w)|^{k+1}} + \frac{|\rho(z)|^{n-k-1}|\rho(w)|^{n-1/2}}{\|z-w\|_A^{2(n-k)}|\Psi(z, w)|^{n+k}}|z-w| \\ = I_1(z, w) + I_2(z, w).$$

Using the coordinate system of Lemma 6.1 with $t = (t_1, t_2, \dots, t_n)$, it follows that

$$\int_{|z-w|<\delta} I_1(z, w) dV(w) \leq C \int_{|t|<1} \frac{|\rho(z)|^{n-k-1}|t_1| dt}{(|t_1|^2 + |\rho(z)||t|^2)^{n-k}(|\rho(z)| + |t|^2)^{k+1}}.$$

Changing to polar coordinates with $s = |t_1|$ and $r = |t'|$, where $t' = (t_2, \dots, t_n)$, we have

(29)

$$\int_{|z-w|<\delta} I_1(z, w) dV(w) \leq C \int_0^1 \int_0^1 \frac{|\rho(z)|^{n-k-1}s^2 ds r^{2n-3} dr}{(s^2 + |\rho(z)|r^2)^{n-k}(|\rho(z)| + r^2)^{k+1}} \\ \leq C \int_0^1 \int_0^1 \frac{|\rho(z)|^{-1/2}s ds r^{2n-3} dr}{(s^2/|\rho(z)| + r^2)^{n-k-1/2}(|\rho(z)| + r^2)^{k+1}} \\ \leq C \int_0^1 \frac{|\rho(z)|^{1/2}r^{2k} dr}{(|\rho(z)| + r^2)^{k+1}} \leq C \int_0^1 \frac{|\rho(z)|^{1/2} dr}{|\rho(z)| + r^2} \leq C.$$

The passage from the next to last line to the last line requires $k \leq n-2$. To estimate the integral of I_2 we use Lemma 6.2(a) and obtain

$$\|z-w\|_A^{2(n-k)} \geq c\|z-w\|_A^{2(n-k)-1}(|G(z, w)|^2 + |\rho(w)||z-w|^2)^{1/2} \\ \geq c\|z-w\|_A^{2(n-k)-1}|\rho(w)|^{1/2}|z-w|.$$

Thus

$$I_2(z, w) \leq \frac{C|\rho(z)|^{n-k-1}}{\|z-w\|_A^{2(n-k)-1}|\Psi(z, w)|^{k+1}}.$$

Using the same coordinate system as before, we obtain

$$\int_{|z-w|<\delta} I_2(z, w) dV(w) \leq C \int_{|t|<1} \frac{|\rho(z)|^{n-k-1} dt}{(|t_1|^2 + |\rho(z)||t|^2)^{n-k-1/2}(|\rho(z)| + |t|^2)^{k+1}} \\ \leq C \int_{|t|<1} \frac{|\rho(z)|^{-1/2} dt}{(|t_1|^2/|\rho(z)| + |t|^2)^{n-k-1/2}(|\rho(z)| + |t|^2)^{k+1}}.$$

After a change to polar coordinates, the last integral is dominated by the second line of (29). Thus $\int_{|z-w|<\delta} I_2(z, w) dV(w) \leq C$. Therefore $v_k(z) \leq C$ for all $z \in D$ and $0 \leq k \leq n-2$. Thus we can write $u_1 = v_1 + \dots + v_{n-2}$, and u_1 is bounded as required.

We set $u_2 = v_{n-1}$ and estimate the gradient of u_2 . Note that $\partial_z G = \partial \rho(z) + \mathcal{E}_1$ and $\bar{\partial}_z G = \mathcal{E}_1$. Therefore

$$|\partial_z(\|z-w\|_A^2) \wedge \partial \rho(z)| + |\bar{\partial}_z(\|z-w\|_A^2) \wedge \bar{\partial} \rho(z)| \leq C(|\rho(z)| + |G|)|z-w|,$$

and

$$|\partial_z \Psi(z, w) \wedge \partial \rho(z)| \leq C(|z-w| + |\Psi(z, w)|).$$

From the expression for L_{n-1} in (27) and the definition of $\|z - w\|_A^2$ it follows that

$$\begin{aligned} & |\partial u_2(z) \wedge \partial \rho(z)| + |\bar{\partial} u_2(z) \wedge \bar{\partial} \rho(z)| \\ & \leq C \int_D \left[\frac{|\rho(z)| + |G| + |z - w|^2}{\|z - w\|_A^2 |\Psi(z, w)|^{2n}} + \frac{(|\rho(z)| + |G|)|z - w|^2}{\|z - w\|_A^2 |\Psi(z, w)|^{2n+1}} \right] |\rho(w)|^{n-1/2} dV(w) \\ & \quad + C \int_D \left[\frac{|\rho(z)|^{1/2} + |z - w|}{\|z - w\|_A^2 |\Psi(z, w)|^{2n}} + \frac{|G| |z - w|}{\|z - w\|_A^2 |\Psi(z, w)|^{2n+1}} \right] |\rho(w)|^n dV(w) \\ & = J_1 + J_2, \end{aligned}$$

where J_1 contains the terms which have neither the factor $|z - w|$ nor $|G|$ in their numerators and J_2 contains the remaining terms.

From previous estimates

$$|G(z, w)| + |\rho(z)| + |\rho(w)| + |z - w|^2 \leq C |\Psi(z, w)|.$$

Using again the coordinate system of Lemma 6.1 we have

$$\begin{aligned} J_1(z) & \leq C \int_D \left[\frac{|\rho(z)| |\rho(w)|^{n-1/2}}{\|z - w\|_A^2 |\Psi(z, w)|^{2n}} + \frac{|\rho(z)|^{1/2} |\rho(w)|^n}{\|z - w\|_A^2 |\Psi(z, w)|^{2n}} \right] dV(w) \\ & \leq C |\rho(z)|^{1/2} \int_D \frac{1}{\|z - w\|_A^2 |\Psi(z, w)|^n} dV(w) \\ & \leq C |\rho(z)|^{1/2} \int_{|t| < 1} \frac{dt}{(|t_1|^2 + |\rho(z)| |t|^2)(|\rho(z)| + |t_1| + |t|^2)^n} \end{aligned}$$

(changing to polar coordinates)

$$\leq C |\rho(z)|^{1/2} \int_0^1 \int_0^1 \frac{r^{2n-3} dr s ds}{(s^2 + |\rho(z)| r^2)(|\rho(z)| + s + r^2)^n}$$

(letting $s = |\rho(z)| s'$, $r = |\rho(z)|^{1/2} r'$)

$$\begin{aligned} & \leq C \int_0^\infty \int_0^\infty \frac{|\rho(z)|^{-1/2} r'^{2n-3} dr' s' ds'}{(s'^2 + r'^2)(1 + s' + r'^2)^n} \\ & \leq C |\rho(z)|^{-1/2} \int_0^\infty \int_0^\infty \frac{r'^{2n-4} dr' ds'}{(1 + s' + r'^2)^n} \\ & \leq C |\rho(z)|^{-1/2} \int_0^\infty \frac{r'^{2n-4} dr'}{(1 + r'^2)^{n-1}} \\ & \leq C |\rho(z)|^{-1/2} \int_0^\infty \frac{dr'}{1 + r'^2} \leq C |\rho(z)|^{-1/2}. \end{aligned}$$

In order to estimate J_2 we use Lemma 6.2, the definition of $\|z - w\|_A^2$ and the inequality $|\rho(z)| + |\rho(w)| + |z - w|^2 \leq C |\Psi(z, w)|$ to get

$$J_2 \leq C \int_D \left(\frac{1}{|z - w| |\Psi(z, w)|^{n+1}} + \frac{1}{|\Psi(z, w)|^{n+3/2}} \right) dV(w).$$

The standard estimate given in [DH, p. 454] now yields $J_2 \leq C |\rho(z)|^{-1/2}$, and we have thus proved that the gradient of u_2 in the complex tangential directions (that is, orthogonal to the gradient of ρ at z) are bounded by $C |\rho(z)|^{-1/2}$.

An application of similar arguments yields that

$$|\partial u_2(z)| + |\bar{\partial} u_2(z)| \leq C/|\rho(z)|,$$

so that the gradient of u_2 in the complex direction parallel to the gradient of ρ behaves like $|\rho(z)|^{-1}$. The lemma is proved. \square

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